Geometrisation

in supersymmetry-invariant cohomology – supergerbes, κ -symmetry, and all that

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Geometrisation of de Rham cocycles (with integer periods) in the form of so-called abelian gerbes has long been known as not only necessary for a rigorous definition of the lagrangean (σ -model) dynamics of charged pointlike particles, strings and branes and its geometric (pre-)quantisation but also exceptionally useful in the canonical description of its symmetries and dualities, classification and field-theoretic realisation of the corresponding defects, as well as in a cohomological description of obstructions against their gauging and classification of the ensuing gauged σ -models.

In the talk, a natural geometrisation scheme for cocycles in the supersymmetry-invariant refinement of the de Rham cohomology on (a class of) homogeneous spaces of supersymmetry Lie supergroups shall be postulated, in close structural analogy with its Graßmann-even ancestor, and illustrated on examples motivated by superstring-theoretic considerations. The geometrisation, based on the classical correspondence between the Cartan-Eilenberg cohomology of the supersymmetry Lie supergroup and the Chevalley-Eilenberg cohomology of its tangent Lie superalgebra, in conjunction with the cohomological description of (equivalence classes of) Lie-superalgebra extensions and a moment-map criterion for their integrability, gives rise to higher-(super)geometric objects termed supergerbes. Various (anticipated) equivariance properties of the supergerbes shall be indicated, including the physically fundamental κ -symmetry, and – time permitting – a higher-geometric variant of the İnönü–Wigner contraction mechanism (for the supersymmetry Lie superalgebras) shall be outlined.

Goal:

Extending the **gerbe-theoretic approach** of the bosonic two-dimensional σ -model to (super-) σ -models with **homogeneous spaces of Lie supergroups** as target spaces, in a manner consistent with **rigid and local supersymmetry**.

Discussion based upon

- **1.** arXiv:1706.05682
- 2. arXiv:1808.04470
- 3. arXiv:1810.00856
- 4. arXiv:1905.05235
- 5. arXiv:1909.xxxxx (in writing)

Part I

Learning from life without spin

Point of departure: The non-linear σ -model

Given a *closed* m_fold Ω_p of dim $\Omega_p = p + 1$ (the worldvolume) & a metric m_fold (M, g) with $\underset{(p+2)}{H} \in Z_{dR}^{p+2}(M)$ (the target space), consider the theory of maps determined by (the PLA for) the **Dirac–Feynman amplitudes**

$$\begin{split} \mathcal{A}_{\mathrm{DF}} &\equiv \exp\left(\frac{\mathrm{i}}{\hbar} \, \mathcal{S}_{\sigma}^{(\mathrm{NG})}[\cdot]\right) \; : \; [\Omega_{\rho}, \mathcal{M}] \longrightarrow \mathrm{U}(1) \\ \mathcal{S}_{\sigma}^{(\mathrm{NG})}[x] &= \int_{\Omega_{\rho}} \sqrt{|\det x^* g|} + \int_{\Omega_{\rho}} \, x^* \left(\mathrm{d}^{-1} \underset{(\rho+2)}{\mathrm{H}}\right), \end{split}$$

describing minimal embeddings deformed by Lorentz-type forces sourced by a Maxwell-type (p + 2)-form field $\underset{(p+2)}{H}$.

Applications: mainly the critical bosonic string (and membrane) theory, but also the effective FT of (certain slow) collective

excitations of spin chains

Problem: May need $[\underset{(\rho+1)}{H}]_{dR} \neq 0$ (*e.g.*, for conformality), and so $\neg \exists_{\substack{B \ (\rho+1)}} \in \Omega^{\rho+1}(M)$: $d\underset{(\rho+1)}{B} = \underset{(\rho+2)}{H}$

E.g., $(M, g) = (G, \kappa_{\mathfrak{g}} \circ (\theta_{L} \otimes \theta_{L})) \implies H_{(3)} = \lambda \kappa_{\mathfrak{g}} \circ (\theta_{L} \wedge \theta_{L} \wedge \theta_{L})$ and the Cartan 3-form $\underset{(3)}{H}$ generates $H^{3}_{dR}(G)$ for G 1-connected

But QM à la Dirac & Feynman requires that we compare amplitudes for cobordant trajectories!

Conclusion: Need $S_{\sigma}^{(NG)}$ with critical points (the EL eqⁿs) as for $[\underset{(\rho+2)}{H}]_{dR} = 0$ but s.t. \mathcal{A}_{DF} is well-defined $\forall x(\Omega_{\rho}) \in \ker \partial_{M}$.

This calls for the use of **Cheeger–Simons differential characters** Hol_{$\mathcal{G}^{(p)}$} \in Hom(Z_{p+1} , U(1)) s.t. Hol_{$\mathcal{G}^{(p)}$} $\circ \partial_{\mathcal{M}}(\cdot) = \exp(\frac{i}{\hbar} \int_{(\cdot)} (H_{p+2})$

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Solution: Fix an arbitrary *good* open cover $\mathcal{O}_M = \{\mathcal{O}_i\}_{i \in I}$ & a tesselation $\triangle_{\Omega_p} = \mathfrak{T}_{p+1} \sqcup \mathfrak{T}_p \sqcup \cdots \sqcup \mathfrak{T}_0$ of Ω_p subordinate to it for $x \in [\Omega_p, M]$, *i.e.*, s.t.

 $\exists_{\iota\in\mathrm{Map}(\triangle_{\Omega_{p}},l)} \forall_{\tau\in\triangle_{\Omega_{p}}} : \mathbf{X}(\tau)\subset\mathcal{O}_{\iota_{\tau}},$

and pull back, along x, a resolution/trivialisation of $\underset{(p+2)}{H}$ over \mathcal{O}_M .

E.g., use
$$\underset{(2)}{b} = (B_i, A_{ij}, g_{ijk}) \in \Omega^2(\mathcal{O}_i) \times \Omega^1(\mathcal{O}_{ij}) \times \mathrm{U}(1)_{\mathcal{O}_{ijk}}$$
 s.t.

 $\underset{(3)}{\mathrm{H}}\!\!\upharpoonright_{\mathcal{O}_i} = \mathsf{d}B_i, \qquad (B_j - B_i)\!\!\upharpoonright_{\mathcal{O}_{ij}} = \mathsf{d}A_{ij}, \qquad (A_{jk} - A_{ik} + A_{ij})\!\!\upharpoonright_{\mathcal{O}_{ijk}} = \mathsf{i}\,\mathsf{d}\log g_{ijk}$

to write (for $X_{\tau} \equiv X \upharpoonright_{\tau}$)

$$S_{\sigma}^{(\text{NG}),\text{top}}[x] = \sum_{p \in \mathfrak{T}_{2}} \left[\int_{p} x_{p}^{*} B_{\iota_{p}} + \sum_{e \in \partial p} \left(\int_{e} x_{e}^{*} A_{\iota_{p}\iota_{e}} - i \sum_{v \in \partial e} \varepsilon_{ev} \log g_{\iota_{p}\iota_{e}\iota_{v}}(x(v)) \right) \right],$$

with \mathcal{A}_{DF} well-defined iff $\delta g_{ijkl} = 1$, so that $Db_{(2)} = (\underset{(3)}{\text{H}} \upharpoonright_{\mathcal{O}_{i}}, 0)$
and $\text{Per}(\underset{(3)}{\text{H}}) \subset 2\pi\mathbb{Z}$ (Dirac's quantisation of charge)

Upshot: As in the Clutching Theorem, the DB (p + 1)-cocycle $\underset{(p+1)}{b}$ geometrises as an abelian bundle *p*-gerbe \mathcal{G}_p of curv $(\mathcal{G}_p) = \underset{(p+2)}{H}$:

$$\pi^*_{YM} \underset{(p+2)}{H} = dE$$

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and $\mathcal{A}_{\mathrm{DF}}^{(\mathrm{NG}),\mathrm{top}}[x] \equiv \mathrm{Hol}_{\mathcal{G}_{p}}(x(\Omega_{p})) = \iota_{p}([x^{*}\mathcal{G}_{p}])$ for a canonical group isomorphism $\iota_{p} : \mathcal{W}^{p+2}(\Omega_{p}; 0) \xrightarrow{\cong} \mathrm{U}(1).$

E.g., an abelian bundle 1-gerbe G_1 [Murray & Stevenson '94-'99]



$$(\mathrm{pr}_2^* - \mathrm{pr}_1^*)\mathrm{B} = \mathrm{curv}(\nabla_L)$$

 $\pi^*_{\mathsf{Y}M^{(3)}} = \mathsf{dB}$

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with the groupoid product μ_L on fibres of L associative.

The origin of Species:



Upshot & spin-off:

• geometric quantisation *via* cohomological transgression [Gawędzki '87, rrS '11]

 $\tau_{\rho} : \mathbb{H}^{\rho+1}(M, \mathcal{D}(\rho+1)^{\bullet}) \longrightarrow \mathbb{H}^{1}(\mathcal{C}_{\rho}M, \mathcal{D}(1)^{\bullet}), \qquad \mathcal{C}_{\rho}M \equiv [\mathcal{C}_{\rho}, M]$

yields a (pre)quantum bundle $\mathcal{H}_{\sigma} = \Gamma_{(\text{pol})}(\mathcal{L}_{\sigma})$, where

 $\mathbb{C}^{\times} \longrightarrow \pi^*_{\mathsf{T}^*\mathcal{C}_{\rho}M}\mathcal{L}_{\mathcal{G}_{\rho}} \otimes \mathcal{I}^{\mathbf{0}}_{\vartheta_{\mathsf{T}^*\mathcal{C}_{\rho}M}} \equiv \mathcal{L}_{\sigma} \,, \, \nabla_{\mathcal{L}_{\sigma}}$

 $\mathsf{P}_{\sigma} \equiv \mathsf{T}^* \mathcal{C}_{\rho} \mathcal{M}, \ \Omega_{\sigma} = \delta \vartheta_{\mathsf{T}^* \mathcal{C}_{\rho} \mathcal{M}} + \pi^*_{\mathsf{T}^* \mathcal{C}_{\rho} \mathcal{M}} \int_{\mathcal{C}_{\rho}} \operatorname{ev}^*_{(\rho+2)} \equiv \operatorname{curv}(\nabla_{\mathcal{L}_{\sigma}})$ and hence, classification of σ -models;

 geometrisation and cohomological classification of duality defects [Fuchs et al. '07, Runkel & rrS '08, rrS '11-'12];

explicit constructions for the 'all' 2d RCFTs via...

... The Universal Gauge Principle (morally, $M \rightarrow M//G_{\sigma}$) [Gawędzki & Reis '02-'03, Gawędzki, Waldorf & rrS '07-'13, rrS '11-'13]

Let G_{σ} be a Lie group with $\text{Lie}(G_{\sigma}) \equiv \mathfrak{g}_{\sigma} = \bigoplus_{A=1}^{N} \langle t_A \rangle$. An action $\lambda_{\cdot} : G_{\sigma} \times M \longrightarrow M$ with the fundamental $\mathcal{K}_A \equiv \mathcal{K}_{t_A} \in \Gamma(TM)$ is

• a global symmetry of the σ -model if

 $\forall_{(g,A)\in \mathcal{G}_{\sigma}\times\overline{1,N}} : \lambda_{g}^{*}g = g \quad \wedge \quad \iota_{\mathcal{K}_{A}}_{(\rho+2)} = -\mathsf{d}_{(\rho)}_{(\rho)} \wedge \quad \lambda_{g}^{*}\mathcal{G}_{\rho} \cong \mathcal{G}_{\rho}.$

can be gauged via the 'minimal coupling' of A ≡ A^A ⊗ t_A ∈ Ω¹(P_G) ⊗ g_σ iff

 [Gawędzki, Waldorf & rrS '07-'13, rrS '11-'13, '19]

 the small gauge anomaly vanishes

$$\iff \left(\bigoplus_{A=1}^{N} C^{\infty}(M, \mathbb{R}) \left(\mathcal{K}_{A, \kappa_{(p)}} \right), \left[\cdot, \cdot \right]_{(p+2)}^{H} \right) \cong \mathfrak{g}_{\sigma} \ltimes_{\lambda} M,$$

where $\left[\cdot, \cdot \right]_{(p+2)}^{H}$ is a $\underset{(p+2)}{H}$ -twisted (à la Ševera–Weinstein)

Vinogradov-type bracket on $\Gamma(TM \oplus_{M,\mathbb{R}} \wedge^{\rho} T^*M)$;

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2. [Gawędzki, Waldorf & rrS '07-'13] the large gauge anomaly vanishes

 \Leftrightarrow \mathcal{G}_{ρ} admits a \mathbf{G}_{σ} -equivariant structure relative to curvature

$$\varrho_{(p+1)} := \sum_{k=1}^{p+1} \frac{(-1)^{p-k}}{k!} \operatorname{pr}_{2}^{*} \alpha_{A_{1}A_{2}\dots A_{k}} \wedge \operatorname{pr}_{1}^{*} \left(\theta_{L}^{A_{1}} \wedge \theta_{L}^{A_{2}} \wedge \dots \wedge \theta_{L}^{A_{k}} \right)$$

on $G_{\sigma} \times M$, with the $\alpha_{A_1A_2...A_k}_{(\rho+1-k)}$ determined by the $(\mathcal{K}_A, \mathcal{K}_A)$, *i.e.*, in particular,

 $\Upsilon_{p} : \lambda^{*}_{\cdot}\mathcal{G}_{p} \xrightarrow{\cong} \mathrm{pr}_{2}^{*}\mathcal{G}_{p} \otimes \mathcal{I}_{(p+1)}^{p}.$

<u>NB</u>: \mathcal{G}_{p} *descends* to $M//G_{\sigma}$ iff $\varrho_{(p+1)} = 0$. We shall call the associated G_{σ} -equivariant structure on \mathcal{G}_{p}

descendable.

The many faces of a G_{σ} -equivariant structure:

- extension of the (p + 1)-cocycle in H^{p+1}(M, D(p + 1)) for G_p to a (p + 1)-cocycle in an extension of the Čech–de Rham bicomplex in the direction of G_σ-cohomology;
- extension of the 0-cell \mathcal{G}_{ρ} to a *descent* $(\rho + 2)$ -tuple $(\mathcal{G}_{\rho}, \Upsilon_{\rho}, \Upsilon_{\rho-1}, \ldots, \Upsilon_{0})$ over $\mathsf{N}^{\bullet}(\mathsf{G}_{\sigma} \ltimes M)$;
- [rrS '12] geometric data for the topological gauge-symmetry defect of the σ -model over Ω_p (based on [Runkel & rrS '09]).

Applications:

- geometrisation and cohomological classification of obstructions against gauging and of inequivalent gaugings, and hence
- natural mapping of the moduli space of σ -models, with beautiful connections to TFT

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reconstruction of T-duality outside the topological context...

Part II

Putting a spin on it or The Brave New Superworld

The goal:

A rigorous definition of a super- σ -model of 'mappings' $[\Omega_p, \mathcal{M}]$ for \mathcal{M} a superm_fold endowed with an action of a supersymmetry Lie supergroup G.

Problems:

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- What is the meaning of (all this, and in particular of) $[\Omega_{\rho}, \mathcal{M}]$?
- Inherent non-compactness of $G \implies H^{\bullet}_{dR}(\mathcal{M})^G \not\equiv H^{\bullet}_{dR}(\mathcal{M})$.

Physical motivation:

Understanding the (super)geometric structure (*sensu largissimo*) of superstring theory-inspired & -related FTs, with view to elucidation of the deep nature of the tremendously robust yet notoriously elusive AdS/CFT correspondence.

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Defⁿ: A **superm_fold** of superdimension (m|n) is a ringed space $\mathcal{M} = (|\mathcal{M}|, \mathcal{O}_{\mathcal{M}})$ composed of

• $|\mathcal{M}| \in \text{Ob} \operatorname{TopMan}$ of dim $|\mathcal{M}| = m$ (the body) and

• $\mathcal{O}_{\mathcal{M}}$: $\mathcal{T}(|\mathcal{M}|)^{\mathrm{op}} \longrightarrow \mathsf{SAlg}_{\mathsf{scomm}}$ (the structure sheaf), locally modelled on $(\mathbb{R}^m, C^{\infty}(\cdot, \mathbb{R}) \otimes \wedge^{\bullet} \mathbb{R}^n) \equiv \mathbb{R}^{m|n}, i.e.,$

 $\exists_{\{\mathcal{U}_i\}_{i\in I}\in \mathrm{Cov}(|\mathcal{M}|)\subset \mathscr{T}(|\mathcal{M}|)} \forall_{j\in I} \exists_{\mathcal{V}_j\in \mathscr{T}(\mathbb{R}^m), \ \mathcal{V}_j\cong \mathcal{U}_j} \ : \ \mathcal{O}_{\mathcal{M}}(\mathcal{U}_j)\cong C^{\infty}(\mathcal{V}_i,\mathbb{R})\otimes \wedge^{\bullet}\mathbb{R}^n \,.$

A superm_fold morphism

 $\varphi \equiv (|\varphi|, \varphi^*) : (|\mathcal{M}_1|, \mathcal{O}_{\mathcal{M}_1}) \longrightarrow (|\mathcal{M}_2|, \mathcal{O}_{\mathcal{M}_2})$

consists of

• $|\varphi| \in \operatorname{Hom}_{\operatorname{TopMan}}(|\mathcal{M}_1|, |\mathcal{M}_2|)$ and

• $\varphi^* : \mathcal{O}_{\mathcal{M}_2} \Longrightarrow |\varphi|_* \mathcal{O}_{\mathcal{M}_1}, i.e.,$ a family $\varphi^*_{\mathcal{U}} \in \operatorname{Hom}_{\mathsf{sAlg}_{\mathsf{scomm}}}(\mathcal{O}_{\mathcal{M}_2}(\mathcal{U}), \mathcal{O}_{\mathcal{M}_1}(|\varphi|^{-1}(\mathcal{U}))$ indexed by $\mathcal{U} \in \mathscr{T}(|\mathcal{M}_2|).$

> Together, these form the **category of superm_folds sMan**, with product $\mathcal{M}_1 \times \mathcal{M}_2 = (|\mathcal{M}_1| \times |\mathcal{M}_2|, \mathcal{O}_{\mathcal{M}_1} \widehat{\otimes} \mathcal{O}_{\mathcal{M}_2}).$

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E.g.,

- $(M, C^{\infty}(\cdot, \mathbb{R}))$ is a superm_fold of superdimension $(\dim M|0)$;
- (M, Ω•(·)) is a superm_fold of superdimension (dim M|dim M);
- super-Minkowski spacetime, super-AdS spacetime, Lie supergroups, their homogeneous spaces...

The geometric perspective: By the Yoneda Lemma,

Yon. : **sMan** \hookrightarrow **Presh**(**sMan**) faithfully, and so we get, for any $(|\mathcal{M}|, \mathcal{O}_{\mathcal{M}}) \in \text{Ob sMan}$,

 $\operatorname{Yon}_{\mathcal{M}}(\cdot) \equiv \operatorname{Hom}_{sMan}(\cdot, \mathcal{M}) : sMan^{\operatorname{op}} \longrightarrow Set$,

with Yon. $(S) \equiv \operatorname{Hom}_{sMan}(S, \mathcal{M})$ the set of S-points in \mathcal{M} . In particular, $\operatorname{Hom}_{sMan}(\mathbb{R}^{0|0}, \mathcal{M}) \equiv |\mathcal{M}|$ is the set of topological points in \mathcal{M} (also, $S \dashrightarrow \mathbb{R}^{0|0} \longrightarrow \mathcal{M}$).

In the S-point picture, we obtain local coordinates

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$$(x^{a}, \theta^{\alpha}), \qquad (a, \alpha) \in \overline{1, m} \times \overline{1, n}$$

of the respective Graßmann parities $|x^a| = 0 = |\theta^{\alpha}| - 1$. Sections take the form

$$f(x,\theta) = \sum_{k=0}^{\infty} F_{\alpha_1\alpha_2...\alpha_k}(x) \otimes \theta^{\alpha_1} \theta^{\alpha_2} \cdots \theta^{\alpha_k}$$

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Furthermore, we have the locally free $\mathbb{Z}/2\mathbb{Z}$ -graded tangent sheaf

 $\mathcal{TM} = \operatorname{sDer} \mathcal{O}_{\mathcal{M}},$

with (homogeneous) sections \mathcal{X} subject to the super-Leibniz rule, written for homogeneous sections f, g of $\mathcal{O}_{\mathcal{M}}$,

$$\mathcal{X}(f \cdot g) = \mathcal{X}(f) \cdot g + (-1)^{|\mathcal{X}| \cdot |f|} f \cdot \mathcal{X}(g),$$

and the dual (locally free $\mathbb{Z}/2\mathbb{Z}$ -graded) cotangent sheaf

 $\mathcal{T}^*\mathcal{M} = \operatorname{Hom}_{Sh}(\mathcal{TM}, \mathcal{O}_{\mathcal{M}}).$

In the local coörds (x^a, θ^{α}) , we obtain the local bases:

 $\left\{\partial_a \equiv \frac{\partial}{\partial x^a}, \vec{\partial}_\alpha \equiv \frac{\vec{\partial}}{\partial \theta^\alpha}\right\}$ of \mathcal{TM} and $\left\{dx^a, d\theta^\alpha\right\}$ of $\bigwedge^{\bullet} \mathcal{T}^* \mathcal{M}$,

and the corresponding presentations of sections:

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Among Ob **sMan**, we find

Defⁿ: A Lie supergroup is a group object

 $(\mathbf{G} = (|\mathbf{G}|, \mathcal{O}_{\mathbf{G}}), \mu : \mathbf{G} \times \mathbf{G} \longrightarrow \mathbf{G}, \operatorname{Inv} : \mathbf{G} \circlearrowleft, \varepsilon : \mathbb{R}^{\mathbf{0}|\mathbf{0}} \longrightarrow \mathbf{G})$

in **sMan**, with body $|G| \in Ob$ LieGrp. Together with (obvious) morphisms, they form the category of Lie supergroups sLieGrp. **NB:** $\forall_{\mathcal{S}\in Ob \ sMan}$: Yon. (\mathcal{S}) : sLieGrp \rightarrow TopGrp. On $G \in Ob \ sLieGrp$, we have LI vector fields $L \in \mathcal{T}G(|G|)$ determined by the condition $\mathcal{O}_G \otimes \mathcal{O}_G \xleftarrow{\mu^*} \mathcal{O}_G$

and their duals - the LI 1-forms.

The RI objects are defined analogously

 $id_{\mathcal{O}_G} \otimes L$

 $\stackrel{\forall}{\mathcal{O}_{G} \widehat{\otimes} \mathcal{O}_{G}} \underbrace{\leftarrow}_{\mu^{*}}$

L,

The supercommutator closes on the LI vector fields, giving rise to the **tangent Lie superalgebra** of G,

 $(\mathrm{sLie}\,(\mathrm{G})\cong\mathsf{T}_{e}\mathrm{G},[\cdot,\cdot\}),\qquad\mathfrak{g}\ni X\longmapsto L_{X}\equiv(\mathrm{id}_{\mathcal{O}_{\mathrm{G}}}\otimes X)\circ\mu^{*}\in\mathcal{T}\mathrm{G}(|\mathrm{G}|)_{\mathrm{L}},$

an example of

Def^{**n**}: A Lie superalgebra is $\mathfrak{g} = \mathfrak{g}^{(0)} \oplus \mathfrak{g}^{(1)} \in Ob$ sVect with the Lie superbracket

$$[\cdot, \cdot\} : \mathfrak{g} \times \mathfrak{g} \longrightarrow \mathfrak{g} \qquad \text{s.t.}$$

(sLB1) $[\mathfrak{g}^{(i)},\mathfrak{g}^{(j)}\} \subset \mathfrak{g}^{(i+2j)};$ (sLB2) $\forall_{X,Y \in \mathfrak{g} \text{ (hom.)}} : [Y,X] = (-1)^{|X| \cdot |Y|} [X,Y];$ (sLB3) the super-Jacobi identity holds true.

Together with the obvious morphisms, they form

the category of Lie superalgebras sLieAlg.

Th^m[Kostant '77]: $\exists \mathcal{K} : sLieGrp \xrightarrow{\cong} sHCp$

Objects and morphisms of the latter category are given in

Def^{\square}: A **super-Harish-Chandra pair** is a... triple $G \equiv (|G|, \mathfrak{g}, \rho)$ composed of

- $|G| \in Ob$ LieGrp,
- $\mathfrak{g} = \mathfrak{g}^{(0)} \oplus \mathfrak{g}^{(1)} \in \operatorname{Ob} \mathbf{sLieAlg}$ with $\mathfrak{g}^{(0)} \equiv \operatorname{Lie}(|G|)$, and
- a realisation $\rho : |\mathbf{G}| \longrightarrow \operatorname{End}_{\mathsf{sLieAlg}}(\mathfrak{g})$ s.t., for any $g \in |\mathbf{G}|$, $\rho(g) \upharpoonright_{\mathfrak{g}^{(0)}} \equiv \mathsf{T}_{e} \operatorname{Ad}_{g}$.

A sHCp morphism (Φ, ϕ) : $(|G_1|, \mathfrak{g}_1, \rho_1) \longrightarrow (|G_2|, \mathfrak{g}_2, \rho_2)$ is composed of

- $\Phi \in \operatorname{Hom}_{\operatorname{LieGrp}}(|G_1|, |G_2|),$
- $\phi \in \operatorname{Hom}_{\mathsf{sLieAlg}}(\mathfrak{g}_1, \mathfrak{g}_2)$ s.t. $\phi \upharpoonright_{\mathfrak{g}^{(0)}} = \mathsf{T}_e \Phi$ and, for any $g \in |\mathsf{G}|, \ \rho_2 \circ \Phi(g) \circ \phi = \phi \circ \rho_1(g).$

<u>Remark</u>: \mathcal{H} uses the Hopf-superalgebra structure on $\mathfrak{U}(\mathfrak{g})$. It yields $sLie(\mathcal{H}(|G|, \mathfrak{g}, \rho)) = \mathfrak{g}$.

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Examples of Lie supergroups:

• $sMink(d, 1 | D_{d,1})$ as an abstract Lie supergroup is

 $\operatorname{sMink}(d, 1 | D_{d,1}) = \left(\mathbb{R}^{d+1}, C^{\infty}(\cdot, \mathbb{R}) \otimes \bigwedge^{\bullet} \mathbb{R}^{D_{d,1}} \right), \qquad D_{d,1} = \dim S_{d,1},$

with $S_{d,1}$ a distinguished Majorana-spinor Cliff ($\mathbb{R}^{d,1}$)-module. It admits global coörds $\{x^a, \theta^\alpha\}^{(a,\alpha)\in\overline{0,d}\times\overline{1,D_{d,1}}}$ and

 $\mu^* : (x^a, \theta^\alpha) \longmapsto (x^a \otimes \mathbf{1} + \mathbf{1} \otimes x^a - \frac{1}{2} \theta^\alpha \otimes (C \Gamma^a)_{\alpha\beta} \theta^\beta, \theta^\alpha \otimes \mathbf{1} + \mathbf{1} \otimes \theta^\alpha),$ $\operatorname{Inv}^* : (x^a, \theta^\alpha) \longmapsto (-x^a, -\theta^\alpha),$

or, equivalently, in the S-point picture,

 $(x_1^a, \theta_1^\alpha) \cdot (x_2^b, \theta_2^\beta) = (x_1^a + x_2^a - \frac{1}{2}\theta_1 \overline{\Gamma}^a \theta_2, \theta_1^\alpha + \theta_2^\alpha), \qquad (x^a, \theta^\alpha)^{-1} = (-x^a, -\theta^\alpha)$

As a sHCp, sMink(d, 1 | $D_{d,1}$) = (Mink(d, 1), smint(d, 1 | $D_{d,1}$) = $\bigoplus_{a=0}^{d} \langle P_a \rangle \oplus \bigoplus_{\alpha=1}^{D_{d,1}} \langle Q_{\alpha} \rangle, 0$),

$$\{Q_{\alpha}, Q_{\beta}\} = \left(C\,\Gamma^{a}\right)_{\alpha\beta}P_{a}\,, \qquad \left[P_{a}, P_{b}\right] = 0 = \left[Q_{\alpha}, P_{a}\right].$$

• SU(2, 2 | 4) as a sHCp with the body Lie group $|SU(2, 2 | 4)| = SO(4, 2) \times SO(6)$,

the Lie superalgebra

$$\mathfrak{su}(2,2|4) = \left(\left(\bigoplus_{a=0}^{4} \langle P_a \rangle \oplus \bigoplus_{a'=5}^{9} \langle P_{a'} \rangle \right) \oplus \bigoplus_{(\alpha,\alpha',l) \in \overline{1,4} \times \overline{1,4} \times \{1,2\}} \langle Q_{\alpha\alpha'l} \rangle \right)$$
$$\oplus \left(\bigoplus_{a,b=0}^{4} \langle J_{ab} = -J_{ba} \rangle \oplus \bigoplus_{a',b'=5}^{9} \langle J_{a'b'} = -J_{b'a'} \rangle \right)$$

 $\{Q_{\alpha\alpha'I}, Q_{\beta\beta'J}\} = i\left(-2(\widehat{C}\,\widehat{\Gamma}^{\widehat{a}}\otimes \mathbf{1})_{\alpha\alpha'I\beta\beta'J}P_{\widehat{a}} + (\widehat{C}\,\widehat{\Gamma}^{\widehat{a}\widehat{b}}\otimes \sigma_2)_{\alpha\alpha'I\beta\beta'J}J_{\widehat{a}\widehat{b}}\right),$

 $[\mathcal{Q}_{\alpha\alpha' I}, \mathcal{P}_{\hat{a}}] = -\frac{1}{2} \left(\widehat{\Gamma}_{\hat{a}} \otimes \sigma_2 \right)_{\alpha\alpha' I}^{\beta\beta' J} \mathcal{Q}_{\beta\beta' J}, \qquad [\mathcal{P}_{\hat{a}}, \mathcal{P}_{\hat{b}}] = \varepsilon_{\hat{a}\hat{b}} J_{\hat{a}\hat{b}}, \quad \varepsilon_{\hat{a}\hat{b}} = \begin{cases} +1 & \text{if } \hat{a}, \hat{b} \in \overline{0, 4} \\ -1 & \text{if } \hat{a}, \hat{b} \in \overline{5, 9} \\ 0 & \text{otherwise} \end{cases}$

 $[J_{\widehat{a}\widehat{b}}, J_{\widehat{c}\widehat{d}}] = \eta_{\widehat{a}\widehat{d}} J_{\widehat{b}\widehat{c}} - \eta_{\widehat{a}\widehat{c}} J_{\widehat{b}\widehat{d}} + \eta_{\widehat{b}\widehat{c}} J_{\widehat{a}\widehat{d}} - \eta_{\widehat{b}\widehat{d}} J_{\widehat{a}\widehat{c}},$

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 $[Q_{\alpha\alpha' I}, J_{\hat{a}\hat{b}}] = -\frac{1}{2} \varepsilon_{\hat{a}\hat{b}} \left(\widehat{\Gamma}_{\hat{a}\hat{b}} \otimes 1 \right)^{\beta\beta' J}_{\alpha\alpha' I} Q_{\beta\beta' J}, \qquad [P_{\hat{a}}, J_{\hat{b}\hat{c}}] = \eta_{\hat{a}\hat{b}} P_{\hat{c}} - \eta_{\hat{a}\hat{c}} P_{\hat{b}}.$

and the standard spinor realisation of the former on the Graßmann-odd component of the latter. Towards homogeneous spaces of Lie supergroups...

Def^{**n**}: A **superm_fold with action** of a Lie supergroup G is a pair (\mathcal{M}, λ) composed of

- $\mathcal{M} \in \operatorname{Ob} \mathbf{sMan}$ and
- $\lambda \in \operatorname{Hom}_{sMan}(G \times \mathcal{M}, \mathcal{M})$ s.t.



Together with the (obvious) morphisms (equivariant supermanifold morphisms), these form the **category of G-superm_folds** G-**sMan**.

Among Ob G-**SMan**, we have G with the L action ℓ . and the R action \wp . induced from μ . These restrict to actions of $|G| \ni g$ by superdiffeo_s *as per*

 $|I|_g \equiv \mu \circ (\widehat{g} \times \mathrm{id}_{\mathrm{G}}) \, : \, \mathbb{R}^{\mathsf{0}|\mathsf{0}} \times \mathrm{G} \cong \mathrm{G} \longrightarrow \mathrm{G} \, ,$

 $r|_g \equiv \mu \circ (\mathrm{id}_{\mathrm{G}} \times \widehat{g}) : \mathrm{G} \times \mathbb{R}^{0|0} \cong \mathrm{G} \longrightarrow \mathrm{G}.$

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Th^m[Kostant '77, Koszul '82, Fioresi *et al.* '07]: Given $G \in Ob \ sLieGrp$ and its closed subsupergroup H with $sLie H \equiv h$, \exists ess. unique superm_fold structure on the homogeneous space

 $G/H = (|G|/|H|, \mathcal{O}_{G/H})$ s.t.

 $\mathcal{O}_{G/H} = \left\{ f \in \mathcal{O}_{G} \upharpoonright_{|G|/|H|} \mid \forall_{(J,h) \in \mathfrak{h} \times |H|} : L_{J}(f) = 0 \land |r|_{h}^{*}(f) = f \right\}$

(sHS1) $\pi_{G/H} = (\pi_{|G|/|H|}, \iota_{\mathcal{O}_{G/H}} : \mathcal{O}_{G/H} \hookrightarrow \mathcal{O}_G)$ is a submersion with

 $G \times G -$

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 $\begin{array}{c|c} \operatorname{id}_{G} \times \pi_{G/H} \\ & & & \\ G \times G/H & \xrightarrow{[\ell]:} & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & &$

(sHS2) ℓ . descends to G/H as per

We have local trivialisations over (distinguished) $\mathcal{U} \in \mathcal{T}(|G|/|H|)$,

 $\tau^{-1} : \widehat{\mathcal{U}} \times \mathrm{H} \equiv \left(\mathcal{U}, \mathcal{O}_{\mathrm{G/H}} \restriction_{\mathcal{U}} \right) \times \mathrm{H} \longrightarrow \left(\pi^{-1}_{|\mathrm{G}|/|\mathrm{H}|}(\mathcal{U}), \mathcal{O}_{\mathrm{G}} \restriction_{\pi^{-1}_{|\mathrm{G}|/|\mathrm{H}|}(\mathcal{U})} \right),$

and so also local sections

$$\sigma_{\widehat{\mathcal{U}}} = \tau^{-1} \circ \left(\operatorname{id}_{\widehat{\mathcal{U}}} \times \widehat{\boldsymbol{e}} \right) \; : \; \widehat{\mathcal{U}} \times \mathbb{R}^{0|0} \cong \widehat{\mathcal{U}} \longrightarrow \mathrm{G} \, .$$

With the help of the $[|I|]_{g_i}$, for some $g_i \in |G|$, $i \in I$, induced from |I|. *via* $[\ell]$, we obtain a trivialising cover from $\widehat{\mathcal{U}}_0 \ni H \equiv |I|_e(H)$,

 $\left\{ [|I|]_{g_i} (\widehat{\mathcal{U}}_0) \equiv \widehat{\mathcal{U}}_i \right\}_{i \in I},$

with the corresponding local trivialising sections

 $\sigma_i = |I|_{g_i} \circ \sigma_{\widehat{\mathcal{U}}_0} \circ [I]_{g_i^{-1}}.$

We may be quite explicit for a class of pairs $(\mathfrak{g}, \mathfrak{h})$...

Reductive homogeneous spaces Let $(G, H \subset |G|)$ with $\mathfrak{t} = \mathfrak{t}^{(0)} \oplus \mathfrak{t}^{(1)} \equiv \bigoplus^{d_0} \langle \mathcal{P}_{\mathcal{a}}
angle \oplus \bigoplus^{d_1} \langle \mathcal{Q}_{lpha}
angle \;, \quad \mathfrak{h} = \bigoplus^{d_S} \langle J_{\kappa}
angle$ $\mathfrak{g} = \mathfrak{t} \oplus \mathfrak{h}$, $\alpha = 1$ a=0be reductive, i.e., $[\mathfrak{h},\mathfrak{t}]\subset\mathfrak{t}$. We may then take $\mathcal{U}_0 \equiv |\widehat{\mathcal{U}}_0|$ sufficiently small to have $\mathcal{U}_0 = \left\{ \exp\left(x^a \otimes P_a\right) \mid x^a : \mathcal{V}_0 \longrightarrow \mathbb{R} \text{ a local chart near } 0 \in \mathbb{R}^{d_0 + 1} \right\}$ and then also, for $\{\theta^{\alpha}\}^{\alpha \in \overline{1,d_1}}$ a global chart on $\mathbb{R}^{0|d_1}$. $\sigma_0(x^a,\theta^\alpha) = \exp\left(\theta^\alpha \otimes Q_\alpha\right) \circ \exp\left(x^a \otimes P_a\right).$ In the S-point picture, at any $\xi \equiv (x, \theta) \in \widehat{\mathcal{U}}_{ij}$, we find $\sigma_i(\xi) = \sigma_i(\xi) \cdot h_{ii}(\xi), \qquad h_{ii}(\xi) \in \mathbf{H}.$

Furthermore, components of the LI g-valued Maurer-Cartan 1-form

$$\theta_{\rm L} = \theta_{\rm L}^{\cal A} \otimes t_{\cal A} = \theta_{\rm L}^{\mu} \otimes T_{\mu} + \theta_{\rm L}^{\kappa} \otimes J_{\kappa} , \qquad \bigoplus_{\mu=0}^{a_0+a_1} \langle T_{\mu} \rangle \equiv \mathfrak{t}$$

• transform as H-tensors for $A = \mu \in \overline{0, d_0 + d_1}$,

 $|\mathbf{r}|^*_{\cdot} heta^{\mu}_{\mathrm{L}}(\boldsymbol{g},\boldsymbol{h}) =
ho(\boldsymbol{h})^{\mu}_{\
u} \, heta^{
u}_{\mathrm{L}}(\boldsymbol{g})\,, \qquad (\boldsymbol{g},\boldsymbol{h}) \in \mathrm{Yon}_{\mathrm{G}} imes \mathrm{Yon}_{\mathrm{H}}$

• compose a principal H-connection on $G \longrightarrow G/H$ as

 $\Theta = \theta_{\mathrm{L}}^{\kappa} \otimes J_{\kappa}$.

Upshot:

• the notion of horizontality ($\in \ker \Theta$);

• $T = \tau_{\mu_1 \mu_2 \dots \mu_n} \theta_L^{\mu_1} \otimes \theta_L^{\mu_2} \otimes \dots \otimes \theta_L^{\mu_n}$ with const. coefficients s.t.

 $\tau_{\mu_1\mu_2...\mu_n} = \tau_{\nu_1\nu_2...\nu_n} \,\rho(h)^{\nu_1}_{\ \mu_1} \,\rho(h)^{\nu_2}_{\ \mu_2} \,\cdots \,\rho(h)^{\nu_n}_{\ \mu_n} \,, \qquad h \in \mathcal{H}$

descend to G/H along the σ_i , $i \in I$.

Examples of reductive homogeneous spaces of Lie supergroups:

• $\operatorname{sMink}(d, 1 | D_{d,1}) \equiv \operatorname{sISO}(d, 1 | D_{d,1}) / \operatorname{SO}(d, 1)$ for $\operatorname{sISO}(d, 1 | D_{d,1}) = \operatorname{sMink}(d, 1 | D_{d,1}) \rtimes_{L_{d,1} \oplus S_{d,1}} \operatorname{SO}(d, 1)$, with

 $\mathbf{g} = \eta_{ab} \, \theta_{\mathrm{L}}^{a} \otimes \theta_{\mathrm{L}}^{b} \,,$

 $\underset{(p+2)}{\mathrm{H}} = \begin{cases} \theta_{\mathrm{L}}^{\alpha} \wedge (\mathcal{C} \, \Gamma_{11})_{\alpha\beta} \, \theta_{\mathrm{L}}^{\beta} & (p=0) \\ \theta_{\mathrm{L}}^{\alpha} \wedge (\mathcal{C} \, \Gamma_{a_{1}a_{2}\dots a_{p}})_{\alpha\beta} \, \theta_{\mathrm{L}}^{\beta} \wedge \theta_{\mathrm{L}}^{a_{1}} \wedge \theta_{\mathrm{L}}^{a_{2}} \wedge \dots \wedge \theta_{\mathrm{L}}^{a_{p}} & (1$

the admissible (d, p, N) filling up the 'old brane scan' • $s(AdS_5 \times S^5) \equiv SU(2, 2|4)/(SO(4, 1) \times SO(5))$, with

$$\begin{split} \mathbf{g} &= \eta_{ab} \,\theta_{\mathrm{L}}^{a} \otimes \theta_{\mathrm{L}}^{b} + \delta_{a'b'} \,\theta_{\mathrm{L}}^{a'} \otimes \theta_{\mathrm{L}}^{b'} \,, \\ \mathbf{H} &= \theta_{\mathrm{L}}^{\alpha\alpha' I} \wedge \big(\widehat{C} \,\widehat{\Gamma}_{\widehat{a}} \otimes \sigma_{3}\big)_{\alpha\alpha' I \,\beta\beta' J} \,\theta_{\mathrm{L}}^{\beta} \wedge \theta_{\mathrm{L}}^{\widehat{a}} \end{split}$$

Defⁿ: Given a *closed* m_fold Ω_p of dim $\Omega_p = p + 1$, a Lie supergroup G and its closed Lie subgroup $H \subset |G|$ with $(\mathfrak{g}, \mathfrak{h})$ reductive, as described above, fix a tesellation $\Delta(\Omega_p)$ of Ω_p subordinate, for a given $\xi \in [\Omega_p, G/H]$, to a trivialising cover $\{\widehat{\mathcal{U}}_i\}_{i \in I}$ of G/H, described earlier. Assume given H-basic LI tensors on G:

$$g = g_{(ab)} \theta^a_L \otimes \theta^b_L, \qquad \qquad \chi_{(p+2)} \equiv \pi^*_{G/H}{}_{(p+2)} \in Z^{p+2}_{dR}(G)^G$$

The Green–Schwarz super- σ -model in the Nambu–Goto formulation is a theory of mappings $\xi \in [\Omega_p, G/H]$ determined by the PLA for the DF amplitudes defined in terms of $(\xi_\tau \equiv \xi \upharpoonright_{\tau})$

$$S_{\mathrm{GS},p}^{(\mathrm{NG})}[\xi] = \sum_{\tau \in \mathfrak{T}_{p+1}} \int_{\tau} \sqrt{\det\left((\sigma_{i_{\tau}} \circ \xi_{\tau})^* \mathrm{g}\right)} + \int_{\Omega_p} \xi^* \mathrm{d}^{-1} \underset{(p+2)}{\mathrm{H}}$$

The meaning of the super- σ -model [Freed '95]: The mappings from

$$[\Omega_{\rho}, \mathcal{M}] \equiv \underline{\mathrm{Hom}}_{\mathsf{sMan}}(\Omega_{\rho}, \mathcal{M}) := \mathrm{Hom}_{\mathsf{sMan}}(\cdot \times \Omega_{\rho}, \mathcal{M})$$

$\in \operatorname{Ob} \operatorname{Fun}(\operatorname{sMan}^{\operatorname{op}}, \operatorname{Set})$

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are to be evaluated on (cp Sorokin's superembedding formalism!)

 $\mathbb{R}^{0\,|\,N}\,,\qquad N\in\mathbb{N}^{\times}$

with the resp. structure sheaves $\mathbb{R}[\eta^1, \eta^2, \dots, \eta^N]$, whereupon

 $\xi^{a} = \xi_{0}^{a} + \xi_{i_{1}i_{2}}^{a} \eta^{i_{1}} \eta^{i_{2}} + \dots + \xi_{i_{1}i_{2}\dots2[\frac{N}{2}]}^{a} \eta^{i_{1}} \eta^{i_{2}} \cdots \eta^{i_{2[\frac{N}{2}]}},$ $\xi^{\alpha} = \xi_{i_{1}}^{\alpha} \eta^{i_{1}} + \xi_{i_{1}i_{2}i_{3}}^{\alpha} \eta^{i_{1}} \eta^{i_{2}} \eta^{i_{3}} + \dots + \xi_{i_{1}i_{2}\dots2[\frac{N-1}{2}]+1}^{\alpha} \eta^{i_{1}} \eta^{i_{2}} \cdots \eta^{i_{2[\frac{N-1}{2}]+1}},$ & the $\xi_{i_{1}i_{2}\dotsi_{k}}^{a}, \xi_{i_{1}i_{2}\dotsi_{k}}^{\alpha}$ become the fields of the super- σ -model.

Physically relevant models:

- (i) the original Green-Schwarz-... super-*p*-branes on $sMink(d, 1 | ND_{d,1}) \equiv sISO(d, 1 | ND_{d,1})/SO(d, 1), N \in \mathbb{N}^{\times};$
- (ii) the Metsaev-Tseytlin superstring and super-*p*-branes on $s(AdS_5 \times S^5) \equiv SU(2,2|4)/(SO(4,1) \times SO(5));$
- (iii) the Zhou super-0-brane and superstring on $s(AdS_2 \times S^2) \equiv SU(1, 1 | 2)_2/(SO(1, 1) \times SO(2));$
- (iv) super-p-branes on s(AdS₃ × S³), s(AdS₄ × S⁷), s(AdS₇ × S⁴),...

Empirical facts:

(H) (i) & (iii) have $\begin{bmatrix} \chi \\ (p+2) \end{bmatrix}_{dR}^{G} = 0$, but $\begin{bmatrix} \chi \\ (p+2) \end{bmatrix}_{dR}^{G} \in CaE^{p+2}(G) \setminus \{0\}$. (iW) (ii) has $\begin{bmatrix} \chi \\ (3) \end{bmatrix}_{dR}^{SU(2,2|4)} = 0 \in CaE^{3}(SU(2,2|4))$, but the supersymmetric primitive does NOT İnönü-Wigner-contract to the one of (i). What are the **PROBLEMS** with the empirical facts?

Ad (IW) Signals apparent ill-definedness of the MT super- σ -model whose construction was based upon the asymptotic correspondence with the GS super- σ -model for sMink(9,1|32). [rrS '18]

Ad (H) The choice of the cohomology critical for the meaning of the topological term in A_{DF} .

AND

Physics (SUSY) favours the Cartan-Eilenberg cohomology

 $\operatorname{CaE}^{\bullet}(G) \equiv H^{\bullet}_{\mathrm{dR}}(G)^{\mathrm{G}}$,

BUT

(How) Does $CaE^{\bullet}(G) \setminus H^{\bullet}(G)$ topologise?

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The Rabin-Crane-type argument/hypothesis:

Secretly, the GS super- σ -model for $[\Omega_{\rho}, G/H \equiv \mathcal{M}]$ is a theory of mappings from $[\Omega_{\rho}, \mathcal{M}/\Gamma_{KR}]$ for $\Gamma_{KR} \subset G$ s.t.

 $\mathcal{M}/\Gamma_{\mathrm{KR}} \cong_{\mathrm{loc.}} \mathcal{M} \wedge H^{\bullet}_{\mathrm{dR}}(\mathcal{M})^{\mathrm{G}} \cong H^{\bullet}_{\mathrm{dR}}(\mathcal{M}/\Gamma_{\mathrm{KR}}).$

A working model:

In the case $\mathcal{M} = sMink(d, 1|D_{d,1})$, the subgroup was identified in [Crane & Rabin '85] as the discrete Kostelecký-Rabin supergroup generated by *integer* supertranslations

$$(\mathbf{x}^{\mathbf{a}}, \theta^{\alpha}) \longmapsto (\mathbf{y}^{\mathbf{b}}, \varepsilon^{\beta}) \cdot (\mathbf{x}^{\mathbf{a}}, \theta^{\alpha})$$

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with $y_{i_1i_2...i_k}^b, \varepsilon_{i_1i_2...i_k}^\alpha \in \mathbb{Z}$ (in the S-point picture).

Field-theoretic consequences: We ought to take into account the Γ_{KR} -twisted sector in $[\Omega_{\rho}, G/H]$, but then the Poisson-Lie superalgebra of the Noether charges of supersymmetry of the GS super- σ -model,

$$\{h_A,h_B\}=f_{AB}^{\ C}h_C+\mathcal{A}_{AB},$$

exhibits a (classical!) wrapping anomaly.

<u>Conclusion</u>: Need to consider extensions of the supersymmetry algebra g.

The latter is merely an (exact) intuition with a rigorous cohomology story behind it...

 Towards geometrisation of supersymmetric de Rham cocycles... Th^m: ∃ isomorphism

 $[\gamma] : H^{\bullet}(\mathfrak{g}, \mathbb{R}) \equiv \mathrm{CE}^{\bullet}(\mathfrak{g}, \mathbb{R}) \xrightarrow{\cong} \mathrm{Ca}\mathrm{E}^{\bullet}(\mathrm{G}) \cong H^{\bullet}_{\mathrm{dR}}(\mathrm{G})^{\mathrm{G}} \,.$

Th^{\underline{m}}: \exists a correspondence

 $CE^{2}(\mathfrak{g},\mathbb{R}) \xleftarrow{1:1} \{ \text{ equivalence classes of supercentral extensions of } \mathfrak{g} \},\$

where



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Th^m: **0** $\longrightarrow \mathbb{R} \longrightarrow \widetilde{\mathfrak{g}}_{[\omega]} \longrightarrow \mathfrak{g} \longrightarrow \mathbf{0}$ determined by $[\omega] \in \operatorname{CaE}^2(G)$ integrates to $\mathbf{1} \longrightarrow \mathbb{C}^{\times} \longrightarrow \widetilde{G} \longrightarrow \mathbf{G} \longrightarrow \mathbf{1}$ iff $\operatorname{Per}(\omega) \subset 2\pi\mathbb{Z}$ and $\ell_{\cdot} : \mathbf{G} \times (\mathbf{G}, \omega) \longrightarrow (\mathbf{G}, \omega)$ has a momentum map.

Idea of geometrisation:

(Inspiration: extended superspacetimes of [de Azcárraga et al.])

1. Look for G-invariant 2-cocycles among the

$$\omega_{\mu_1\mu_2\dots\mu_p} := \iota_{L_{\mu_1}} \iota_{L_{\mu_2}} \cdots \iota_{L_{\mu_p}} \chi_{(p+2)}$$

- 2. Associate a supercentral extension $\mathbb{C}^{\times} \longrightarrow \widetilde{G} \xrightarrow{\widetilde{\pi}} G$ to it and partially reduce $\widetilde{\pi}^* \chi_{(\rho+2)}$ in CaE•(\widetilde{G}).
- 3. Repeat 1.-2. until complete reduction of $\hat{\pi}^* \underset{(p+2)}{\chi}$ is obtained over

an extension $\widehat{G} \xrightarrow{\widehat{\pi}} G$ in the corresponding $CaE^{\bullet}(\widehat{G})$, *i.e.*,

$$\exists \underset{(p+1)}{\beta} \in \Omega^{p+1}(\widehat{\mathrm{G}})^{\widehat{\mathrm{G}}} : \mathsf{d}_{(p+1)}^{\beta} = \widehat{\pi}^*_{(p+2)}^* \chi$$

4. Descend $\beta_{(p+1)}$ to \widehat{G}/H .

5. Use \widehat{G}/H as THE surjective submersion of \mathcal{G}_p & continue à la [Murray & Stevenson *et al.*].

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Constructive results:

Theorem I [rrS '17('12)] Consecutive resolution, through central extensions, of the various CaE super-2-cocycles encountered in the analysis of the GS super-(p + 2)-cocycles on sMink $(d, 1 | (N \cdot)D_{d,1})$, induces a hierarchy of surjective submersions necessary for the geometrisation of the latter, leading to the emergence of the corresponding Green–Schwarz super-*p*-gerbes (explicited for $p \in \{0, 1, 2\}$).

Abstraction: The super-Minkowskian construction yields **Cartan-Eilenberg super-p-gerbes** that are, *morally*, *p*-gerbe objects in **sLieGrp**. (sMink($d, 1 | (N \cdot)D_{d,1}$) is a Lie supergroup!) More generally, we obtain *p*-gerbes with a lift of [ℓ]. to the constitutive surjective submersions that preserves the entire

connective structure.

<u>Constructive results – $ct^{\underline{d}}$:</u>

• The success of the super-Minkowskian geometrisation (of Theorem I) is repeated in [rrS '18] in the setting of Zhou's super- σ -model of [Zhou '99] for the superparticle in $s(AdS_2 \times S^2)$.

• The celebrated Metsaev–Tseytlin super- σ -model of [Metsaev & Tseytlin '98] for the superstring in s(AdS₅ × S⁵), on the other hand, is problematic. There exists an Inönü–Wigner-*non*contractible trivial super-1-gerbe, and a collection of no-go theorems.

Constructive results – $ct^{\underline{d}}$:

Theorem II [rrS '19('17)] The GS super-*p*-gerbes of Theorem I with $p \in \{0, 1\}$ are endowed with a canonical **supersymmetric** Ad.-equivariant structure.

NB: This conforms with the purely even (WZW) story.

But there is even more verifiable Physics in the construction...

A self-consistent dscription of the supersymmetric vacuum

The translational component of the global supersymmetry G of $\mathcal{A}_{DF}^{GS,p(NG)}$ is broken spontaneously by the (class.) vacuum Ψ_{VAC} . \implies Also its Graßmann-odd component has to be reduced on Ψ_{VAC} .

Q: How to remove the spurious (Goldstone) Graßmann-odd fields?

<u>A</u>: A κ -symmetry extension of \mathfrak{h} on Ψ_{VAC} , spanning the latter as its (super)diffeo_s. [de Azcárraga & Lukierski '82, Siegel '83]

Problems: Ω_p -locality of κ -symmetry & the mixing of the metric and topological DOFs prevent geometrisation in the NG picture

BUT

Here come Hughes & Polchinski...

Consider $\mathfrak{t}_{VAC}^{(0)} \subset \mathfrak{t}_{VAC}^{(0)} \oplus \mathfrak{e}^{(0)} \subset \mathfrak{t}^{(0)} (\subset \mathfrak{t} \subset \mathfrak{t} \oplus \mathfrak{h} \equiv \mathfrak{g})$ representing Ψ_{VAC} , & its maximal ad-isotropy algebra $\mathfrak{h}_{VAC} \subset \mathfrak{h}$ over $H_{VAC} \subset H$. Write

$$\mathfrak{t} \oplus \mathfrak{h} = \mathfrak{g} = (\mathfrak{t} \oplus \mathfrak{d}) \oplus \mathfrak{h}_{\mathrm{VAC}} \equiv \mathfrak{f} \oplus \mathfrak{h}_{\mathrm{VAC}}, \qquad \mathfrak{d} = \bigoplus \left\langle J_{\widehat{\mathbf{S}}} \right\rangle$$

 $\hat{S}=1$

& assume reductivity of the latter, $[\mathfrak{h}_{VAC}, \mathfrak{f}] \subset \mathfrak{f}$, and $[\mathfrak{h}_{VAC}, \mathfrak{e}^{(0)}] \subset \mathfrak{e}^{(0)}$, $[\mathfrak{d}, \mathfrak{t}_{VAC}^{(0)}] \subset \mathfrak{e}^{(0)}$, $[\mathfrak{d}, \mathfrak{e}^{(0)}] \subset \mathfrak{t}_{VAC}^{(0)}$, as well as unimodularity of ρ on Ψ_{VAC} (volume preservation)

$$\forall h \in \mathrm{H}_{\mathrm{VAC}} : \det \rho(h) \upharpoonright_{\mathfrak{t}_{\mathrm{VAC}}^{(0)}} \equiv \det \mathsf{T}_{e} \mathrm{Ad}_{h} \upharpoonright_{\mathfrak{t}_{\mathrm{VAC}}^{(0)}} \stackrel{!}{=} 1.$$

Replace the former local sections of $G \longrightarrow G/H$ by those of $G \longrightarrow G/H_{VAC}$ given by top. translates of

$$\sigma_{0}^{\text{VAC}}(x^{a}, \theta^{\alpha}, \phi^{\widehat{S}}) = \exp(\theta^{\alpha} \otimes Q_{\alpha}) \cdot \exp(x^{a} \otimes P_{a}) \cdot \exp(\phi^{\widehat{S}} \otimes J_{\widehat{S}})$$
$$\equiv \sigma_{0}(x^{a}, \theta^{\alpha}) \cdot \exp(\phi^{\widehat{S}} \otimes J_{\widehat{S}})$$

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Th^m[rrS '19('17)]: If $(\mathfrak{g}, \mathfrak{h}, \mathfrak{h}_{VAC}, \mathfrak{t}_{VAC}^{(0)})$ and ρ are constrained as above & there exists a T_eAd_H-invariant scalar product g on $\mathfrak{t}^{(0)}$ s.t. $\mathfrak{t}_{VAC}^{(0)} \perp_{\mathfrak{g}} \mathfrak{e}^{(0)}$

then $\mathcal{A}_{DF}^{GS,p(NG)}$ in the gauge σ_0 is (classically) equivalent to the **Green–Schwarz super-\sigma-model in the Hughes–Polchinski** formulation for $[\Omega_p, G/H_{VAC}] \ni \hat{\xi}$ in the gauge σ_0^{VAC} with

$$S_{\mathrm{GS},\rho}^{(\mathrm{HP})}[\widehat{\xi}] = \sum_{\tau \in \mathfrak{T}_{p+1}} \int_{\tau} \left(\sigma_{i_{\tau}}^{\mathrm{VAC}} \circ \widehat{\xi} \right)^* \left(\beta_{(p+1)}^{(\mathrm{HP})} + \mathsf{d}^{-1} \pi_{\mathrm{G/H}(p+2)}^* \right),$$

$$\beta_{(p+1)}^{(\mathrm{HP})} = \frac{\lambda_p}{(p+1)!} \epsilon_{\underline{a}_0 \underline{a}_1 \dots \underline{a}_p} \theta_{\mathrm{L}}^{\underline{a}_0} \wedge \theta_{\mathrm{L}}^{\underline{a}_1} \wedge \dots \wedge \theta_{\mathrm{L}}^{\underline{a}_p}, \qquad \lambda_p \in \mathbb{R}^{\times},$$

perticults reduced through imposition the Lemma High Control

partially reduced through imposition the Inverse Higgs Constraints

 $\left(\sigma_{i_{\tau}}^{\mathrm{VAC}}\circ\widehat{\xi}\right)^{*}\mathsf{P}_{\mathfrak{e}^{(0)}}^{\mathfrak{g}}\circ\theta_{\mathrm{L}}\stackrel{!}{=}\mathsf{0}\,,$

 \iff the EL eqⁿs for the Goldstone modes ϕ^{S} are imposed.

Inspiration: [Hughes & Polchinski '86; Gauntlett et al. '90].

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For the reduction of the Graßmann-odd DOFs, assume

 $\exists \mathsf{P}_{VAC}^{(1)} = \mathsf{P}_{VAC}^{(1)} \circ \mathsf{P}_{VAC}^{(1)} \text{ w/ im } \mathsf{P}_{VAC}^{(1)} \equiv \mathfrak{t}_{VAC}^{(1)} \subset \mathfrak{t}_{VAC}^{(1)} \oplus \mathfrak{e}^{(1)} \equiv \mathfrak{t}^{(1)} \subset \mathfrak{t} \quad \text{s.t.}$ $[\mathfrak{h}_{VAC}, \mathfrak{t}_{VAC}^{(1)}] \subset \mathfrak{t}_{VAC}^{(1)}, \quad [\mathfrak{h}_{VAC}, \mathfrak{e}^{(1)}] \subset \mathfrak{e}^{(1)} \quad \wedge \quad \{\mathfrak{t}_{VAC}^{(1)}, \mathfrak{t}_{VAC}^{(1)}\} \subset \mathfrak{t}_{VAC}^{(0)} \oplus \mathfrak{h}$ & the gauge invariance

 $\forall_{\kappa \in [\Omega_{p}, \mathfrak{t}_{VAC}^{(1)}]} : \frac{d}{dt} \upharpoonright_{t=0} \mathcal{A}_{DF}^{GS, p(HP)} \left[\Phi_{L_{\kappa}} \left(\widehat{\xi}(\cdot), t \right) \right] = 0$ realised *exactly* as $\iota_{L_{\kappa}} \left(\mathsf{d}_{(p+1)}^{\beta(HP)} + \pi_{G/H_{(p+2)}}^{*} \right) = 0$. Then*, the **vacuum slice** Σ_{VAC} within

 $\Sigma = \bigsqcup_{i \in I} \sigma_i^{\text{VAC}}(\widehat{\mathcal{U}}_i^{\text{VAC}}), \qquad \left\{ [|I|]_{g_i}^{\text{VAC}}(\widehat{\mathcal{U}}_0^{\text{VAC}}) \equiv \widehat{\mathcal{U}}_i^{\text{VAC}} \right\}_{i \in I} \in \text{cov}(G/H_{\text{VAC}})$ with the normal determined by the EL eqⁿs of $\mathcal{A}_{\text{DF}}^{\text{GS}, p(\text{HP})}$:

 $\mathsf{P}^{\mathfrak{g}}_{\mathfrak{e}^{(0)}\oplus\mathfrak{e}^{(1)}\oplus\mathfrak{d}}\circ\theta_{\mathsf{L}}\upharpoonright_{\mathcal{T}\Sigma_{\mathsf{VAC}}}=0$ is covered by local flows of **linearised local** κ -symmetries: $\left\langle \mathcal{K}^{\mathsf{R}}_{\delta\xi} \equiv \delta\xi^{\mu}\otimes\left(L_{t_{\mu}}+\Delta^{\underline{S}}_{i_{\mu}}L_{J_{\underline{S}}}\right)\right\rangle_{\delta\xi^{\mu}\in\mathsf{E}_{\mathsf{corf}}(\mathcal{O}_{\mathsf{core}}=0)}$

Problems with the symmetries of the vacuum:

- they do not span an involutive distribution ⇐ their Lie superbracket closes (over Σ_{VAC}!) only up to 'invisible' gauge transformations from h;
- they do not descend to G/H_{VAC} .

Way out: κ -symmetry as a **gauged supervector-space symmetry of the vacuum slice** Σ_{VAC} .

Reminder: Gauge symmetries of the (super-) σ -model geometrise as equivariant structures on the corresponding (super-)p-gerbes, of which the flat ones descend the models and the gerbes to the orbispace.

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Upshot: Need to check the existence of a **supersymmetric linearised descendable** κ -equivariant structure on...

Remarkable upshot: The extended Hughes–Polchinski p-gerbe

$$\widehat{\mathcal{G}}_{\mathrm{HP}}^{(oldsymbol{p})} := \pi^*_{\mathrm{G/H}} \mathcal{G}_{\mathrm{GS}}^{(oldsymbol{p})} \otimes \mathcal{I}^{oldsymbol{p}}_{\begin{subarray}{c} eta^{(\mathrm{HP})} \ (oldsymbol{p}+1) \ (oldsymbol{p}+1) \ \end{array}},$$

written in terms of $\pi_{G/H}$: $G \longrightarrow G/H$, to be restricted to Σ_{VAC} .

<u>NB</u>: The metric DOFs of the NG formulation have been topologised/gerbified!

Theorem III [rrS '19, in writing] $\widehat{\mathcal{G}}_{HP}^{(p)}$ carries a *canonical* linearised (descendable) κ -equivariant structure

 $\Upsilon_{p}^{\kappa}: \mathscr{L}_{\mathcal{K}^{\mathbb{P}}} \,\widehat{\mathcal{G}}_{\mathrm{HP}}^{(p)} \cong \mathcal{I}_{0}^{p} \quad \text{over} \quad \mathfrak{g}_{\mathrm{VAC}} \times \Sigma_{\mathrm{VAC}}.$

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Q: What about compatibility with the global supersymmetry?

Problem: Large supersymmetry transformations $g \in G$ map between sections in Σ_{VAC} , & κ -symmetries do NOT glue...

Way out: Consider linearised supersymmetry,

 $\Phi_{\mathcal{A}} : \mathscr{L}_{\mathcal{K}^{\ell}_{\mathcal{A}}} \widehat{\mathcal{G}}^{(p)}_{\mathrm{HP}} \cong \mathcal{I}^{p}_{0}, \quad \mathcal{A} \in \overline{1, \dim \mathfrak{g}} \quad \text{over} \quad \Sigma_{\mathrm{VAC}},$ implying – in the light of Th^{<u>m</u>} III –

 $\mathscr{L}_{\widetilde{\mathcal{K}}_{A}^{\ell}}\Upsilon_{p} : \mathscr{L}_{\widetilde{\mathcal{K}}_{A}^{\ell}}\mathscr{L}_{\mathcal{K}}^{\wp} \widehat{\mathcal{G}}_{HP}^{(p)} \cong \mathcal{I}_{0}^{p}, \quad A \in \overline{1, \dim \mathfrak{g}} \quad \text{over} \quad \mathfrak{g}_{VAC} \times \Sigma_{VAC}$ for natural lifts $\widetilde{\mathcal{K}}_{A}^{\ell}$ of the \mathcal{K}_{A}^{ℓ} . We demand the existence of



Theorem IV [rrS '19, in writing] The $\{\widetilde{\mathcal{K}}_{A}^{\ell}\}_{A \in \overline{1, \dim \mathfrak{g}}}$ exist, and so Υ_{D}^{κ} is *canonically* supersymmetric (the γ_{A}^{κ} exist).

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Conclusions:

- The physically relevant CaE super-(p + 2)-cocycles on (supersymmetry) Lie supergroups should be geometrised and do geometrise in a large class of supergeometric settings as the GS super-p-gerbes of [rrS '17].
- 2. So do the supersymmetries, global and local. [rrS '19, in writing]
- 3. The super-*p*-gerbes are endowed with (the expected and) natural equivariant structures with respect to the supersymmetries of the relevant super- σ -models, in conformity with the underlying physics and the bosonic intuition. [rrS '19]
- 4. The construction generalises to physically relevant curved homogeneous spaces of supersymmetry Lie supergroups, and sometimes suggests corrections to the existing field-theory results. [rrS '18]

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Outlook:

- Uniqueness of the construction and its relation to the approach of Schreiber *et al.*. Reconstruction of the full-fledged (weak) (p + 1)-categories of super-p-gerbes.
- The relevance of the İnönü–Wigner-contractibility & the ultimate fate of the Metsaev–Tseytlin background.
- The higher supergeometry and superalgebra of supersymmetric defects (including boundary states) & state fusion.
- Relation to the worldvolume supersymmetry, possibly *via* Sorokin's Superembedding Formalism.
- Relation to the String-structure.
- The bosonisation/fermionisation defect.
- T-duality *via* the Hughes–Polchinski formulation, also in the bosonic setting.

• The gauging of the Ad.-supersymmetry and the ensuing hern–Simons-type sTFT. Un (super)buchet pentru Organizatori – din Warszawa, orașul lui Samuel Eilenberg, cu reconuștința și simpatie!



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Dac

Part III

Xtras

Intrusion: The Lie supergroup of the Metsaev-Tseytlin super- σ -model: SU(2,2|4) with the body $|SU(2,2|4)| = SO(4,2) \times SO(6)$

and the Lie superalgebra (*R*-rescaled, for $R \in \mathbb{R}$)

$$\mathfrak{su}(2,2|4)^{(R)} = \left(\left(\bigoplus_{a=0}^{4} \langle P_a \rangle \oplus \bigoplus_{a'=5}^{9} \langle P_{a'} \rangle \right) \oplus \bigoplus_{(\alpha,\alpha',l)\in\overline{1,4\times\overline{1,4\times\{1,2\}}}} \langle Q_{\alpha\alpha'l} \rangle \right) \\ \oplus \left(\bigoplus_{a,b=0}^{4} \langle J_{ab} = -J_{ba} \rangle \oplus \bigoplus_{a',b'=5}^{9} \langle J_{a'b'} = -J_{b'a'} \rangle \right)$$

 $\{Q_{\alpha\alpha' I}, Q_{\beta\beta' J}\} = \mathrm{i}\left(-2(\widehat{C}\,\widehat{\Gamma}^{\widehat{a}}\otimes \mathbf{1})_{\alpha\alpha' I\beta\beta' J}P_{\widehat{a}} + \frac{1}{R^2}\,(\widehat{C}\,\widehat{\Gamma}^{\widehat{a}\widehat{b}}\otimes\sigma_2)_{\alpha\alpha' I\beta\beta' J}J_{\widehat{a}\widehat{b}}\right),$

 $\begin{bmatrix} Q_{\alpha\alpha'I}, P_{\hat{a}} \end{bmatrix} = -\frac{1}{2R} \left(\widehat{\Gamma}_{\hat{a}} \otimes \sigma_2 \right)^{\beta\beta'J}_{\alpha\alpha'I} Q_{\beta\beta'J}, \qquad \begin{bmatrix} P_{\hat{a}}, P_{\hat{b}} \end{bmatrix} = \frac{1}{R^2} \varepsilon_{\hat{a}\hat{b}} J_{\hat{a}\hat{b}}, \quad \varepsilon_{\hat{a}\hat{b}} = \begin{cases} +1 & \text{if } \hat{a}, \hat{b} \in \overline{0, 4} \\ -1 & \text{if } \hat{a}, \hat{b} \in \overline{5, 9} \\ 0 & \text{otherwise} \end{cases},$

 $[J_{\widehat{a}\widehat{b}}, J_{\widehat{c}\widehat{d}}] = \eta_{\widehat{a}\widehat{d}} J_{\widehat{b}\widehat{c}} - \eta_{\widehat{a}\widehat{c}} J_{\widehat{b}\widehat{d}} + \eta_{\widehat{b}\widehat{c}} J_{\widehat{a}\widehat{d}} - \eta_{\widehat{b}\widehat{d}} J_{\widehat{a}\widehat{c}},$

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$$[Q_{\alpha\alpha' I}, J_{\widehat{ab}}] = -\frac{1}{2} \varepsilon_{\widehat{ab}} \left(\widehat{F}_{\widehat{ab}} \otimes \mathbf{1} \right)^{\beta\beta' J}_{\alpha\alpha' I} Q_{\beta\beta' J}, \qquad [P_{\widehat{a}}, J_{\widehat{bc}}] = \eta_{\widehat{ab}} P_{\widehat{c}} - \eta_{\widehat{ac}} P_{\widehat{b}}.$$

with the **inönü-Wigner** asymptote $\mathfrak{su}(2,2|4)^{(R)} \xrightarrow{R \to \infty} \mathfrak{smin}\mathfrak{k}(9,1|32)$

Some Lie-superalgebra cohomology...

Def^{**n**}: A (**left**) $\hat{\mathfrak{g}}$ -module of an LSA $\hat{\mathfrak{g}}$ is a pair (\hat{V}, ℓ .) composed of a **K**-linear superspace $\hat{V} = \hat{V}^{(0)} \oplus \hat{V}^{(1)}$ and a left $\hat{\mathfrak{g}}$ -action

$$\ell_{\cdot} : \widehat{\mathfrak{g}} \times \widehat{V} \longrightarrow \widehat{V} : (X, v) \longmapsto X \triangleright v$$

consistent with the $\mathbb{Z}/2\mathbb{Z}$ -gradings, $X \triangleright v = \tilde{X} + \tilde{v}$, and such that for any two homogeneous elements $X_1, X_2 \in \mathfrak{g}$ and $v \in \hat{V}$,

 $[X_1, X_2\} \vartriangleright v = X_1 \vartriangleright (X_2 \vartriangleright v) - (-1)^{\widetilde{X_1} \cdot \widetilde{X_2}} X_2 \vartriangleright (X_1 \vartriangleright v).$

and the fundamental...

Def^{**n**}: Let $(\hat{\mathfrak{g}}, [\cdot, \cdot])$ be an LSA over field \mathbb{K} and let (\hat{V}, ℓ) be a $\hat{\mathfrak{g}}$ -module. A *p*-cochain on $\hat{\mathfrak{g}}$ with values in \hat{V} is a *p*-linear map $\varphi : \hat{\mathfrak{g}}^{\times p} \longrightarrow \hat{V}$ that is totally super-skewsymmetric,

$$\begin{array}{l} \varphi(X_{1}, X_{2}, \ldots, X_{i-1}, X_{i+1}, X_{i}, X_{i+2}, X_{i+3}, \ldots, X_{p}) \\ = & -(-1)^{\widetilde{X}_{i}\widetilde{X_{i+1}}} \varphi(X_{1}, X_{2}, \ldots, X_{p}) \,. \end{array}$$

They form a \mathbb{Z}_2 -graded group of ρ -cochains on $\widehat{\mathfrak{g}}$ valued in \widehat{V} ,

$$C^p(\widehat{\mathfrak{g}},\widehat{V})=C^p_0(\widehat{\mathfrak{g}},\widehat{V})\oplus C^p_1(\widehat{\mathfrak{g}},\widehat{V})\,,$$

with $\varphi(X_1, X_2, ..., X_p) \in \widehat{V}_{\sum_{i=1}^p \widetilde{X}_i + n}$ for $\varphi \in C_n^p(\widehat{\mathfrak{g}}, \widehat{V})$, composed of even (n = 0) and odd (n = 1) *p*-cochains.

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These groups form a semi-bounded complex $C^{\bullet}(\widehat{\mathfrak{g}}, \widehat{V}) : C^{0}(\widehat{\mathfrak{g}}, \widehat{V}) \xrightarrow{\delta_{\widehat{\mathfrak{g}}}^{(0)}} C^{1}(\widehat{\mathfrak{g}}, \widehat{V}) \xrightarrow{\delta_{\widehat{\mathfrak{g}}}^{(1)}} \cdots \xrightarrow{\delta_{\widehat{\mathfrak{g}}}^{(p-1)}} C^{p}(\widehat{\mathfrak{g}}, \widehat{V}) \xrightarrow{\delta_{\widehat{\mathfrak{g}}}^{(p)}} \cdots$

e as

$$\begin{split} \delta_{\mathfrak{g}}^{(p)} &: C_{n}^{p}(\mathfrak{g}, V) \longrightarrow C_{n}^{p+1}(\mathfrak{g}, V) \\ \text{valuate on the homogeneous } X_{i} \in \mathfrak{g}, \ i \in \overline{0, p+1}, \ \varphi \in C^{p}(\mathfrak{g}, V) \\ (\delta_{\mathfrak{g}}^{(0)} \varphi)(X) &:= (-1)^{|X_{0}| \cdot |\varphi|} X_{0} \rhd \varphi, \\ (\delta_{\mathfrak{g}}^{(0)} \varphi)(X) &:= (-1)^{|X_{0}| \cdot |\varphi|} X_{0} \rhd \varphi, \\ (\delta_{\mathfrak{g}}^{(p)} \varphi)(X_{1}, X_{2}, \dots, X_{p+1}) &:= \sum_{j=1}^{p+1} (-1)^{|X_{j}| \cdot |\varphi| + S(X_{j})} X_{j} \rhd \varphi(X_{1}, X_{2}, \dots, X_{p+1}) \\ &+ \sum_{1 \leq j < k \leq p+1} (-1)^{S(X_{j}) + S(X_{k}) + |X_{j}| \cdot |X_{k}|} \varphi([X_{j}, X_{k}], X_{1}, X_{2}, \dots, X_{p+1}), \\ S(X_{i}) &:= |X_{i}| \cdot \sum_{j=1}^{i-1} |X_{j}| + i - 1. \end{split}$$

The $\mathbb{Z}/2\mathbb{Z}$ -graded V-valued cohomology groups of g are $H^{p}(\mathfrak{g}, V) := H^{p}_{0}(\mathfrak{g}, V) \oplus H^{p}_{1}(\mathfrak{g}, V), \quad H^{p}_{n}(\mathfrak{g}, V) := \frac{\ker \delta^{(p)}_{\mathfrak{g}} |_{C^{p}_{n}(\mathfrak{g}, V)}}{\operatorname{im} \delta^{(p-1)}_{\mathfrak{g}} |_{C^{p-1}_{n}(\mathfrak{g}, V)}}.$

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Def^{**n**}: Let $(\hat{\mathfrak{g}}, [\cdot, \cdot]_{\hat{\mathfrak{g}}})$ and $(\hat{\mathfrak{a}}, [\cdot, \cdot]_{\hat{\mathfrak{a}}})$ be LSAs over field \mathbb{K} . A **supercentral extension of** $\hat{\mathfrak{g}}$ by $\hat{\mathfrak{a}}$ is an LSA $(\tilde{\mathfrak{g}}, [\cdot, \cdot]_{\tilde{\mathfrak{g}}})$ over \mathbb{K} that determines a short exact sequence of LSAs

$$\mathbf{0} \longrightarrow \mathfrak{a} \xrightarrow{\mathcal{I}_{\widehat{\mathfrak{a}}}} \widetilde{\mathfrak{g}} \xrightarrow{\pi_{\widehat{\mathfrak{g}}}} \mathfrak{g} \longrightarrow \mathbf{0} \,,$$

written in terms of an LSA mono $j_{\widehat{\mathfrak{a}}}$ and of an LSA epi $\pi_{\widehat{\mathfrak{g}}}$, and s.t. $j_{\widehat{\mathfrak{a}}}(\widehat{\mathfrak{a}}) \subset \mathfrak{z}(\widetilde{\mathfrak{g}})$ (the supercentre of $\widetilde{\mathfrak{g}}$). Whenever $\pi_{\mathfrak{g}}$ admits an LSA section, *i.e.*, there exists

$$\sigma \in \operatorname{Hom}_{{
m sLie}}(\widehat{\mathfrak{g}}, \widetilde{\mathfrak{g}}), \qquad \pi_{\widehat{\mathfrak{g}}} \circ \sigma = \operatorname{id}_{\widehat{\mathfrak{g}}},$$

the supercentral extension is said to split.

An equivalence of supercentral extensions $\tilde{\mathfrak{g}}_{\alpha}, \alpha \in \{1, 2\}$ of $\hat{\mathfrak{g}}$ by $\hat{\mathfrak{a}}$ is represented by a commutative diagram of LSAs

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The relevant one-way ticket:

Given an LSA $(\mathfrak{g}, [\cdot, \cdot]_{\mathfrak{g}})$ and its supercommutative module \mathfrak{a} , as well as a representative $\Theta \in Z_0^2(\mathfrak{g}, \mathfrak{a})$ of a class in $H_0^2(\mathfrak{g}, \mathfrak{a})$, we define

 $\widetilde{\mathfrak{g}}:=\mathfrak{a}\oplus\mathfrak{g}$

and put on it the Lie superbracket

$$[\cdot,\cdot]_{\Theta} \quad : \quad \widetilde{\mathfrak{g}} \times \widetilde{\mathfrak{g}} \longrightarrow \widetilde{\mathfrak{g}}$$

: $((A_1, X_1), (A_2, X_2)) \longmapsto (\Theta(X_1, X_2), [X_1, X_2]_g).$