

ALGEBROIDS

- A PLAYGROUND FOR GEOMETRIC MECHANICS



BIESZCZADY MOUNTAINS, ON OUR WAY TO
KRZEMIENIEC / КРЕМЕНЕЦЬ / KREMENEC / KREMENÁROS

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INTRODUCTION

HAMILTONIAN MECHANICS

- ASSOCIATED WITH A PHASE SPACE \mathcal{P} WHICH IS SYMPLECTIC (\mathcal{P}, ω) OR POISSON (\mathcal{P}, Λ) MANIFOLD
- ω OR Λ ARE USED TO PRODUCE A VECTOR FIELD FROM A HAMILTONIAN FUNCTION $\mathcal{H}: \mathcal{P} \rightarrow \mathbb{R}$

SYMPLECTIC $X_H \lrcorner \omega = d\mathcal{H}$

POISSON $d\mathcal{H} \lrcorner \Lambda = X_H$

IN DARBOUX COORDINATES $\dot{q}^i = \frac{\partial \mathcal{H}}{\partial p_i} \quad \dot{p}^j = -\frac{\partial \mathcal{H}}{\partial q^j}$

USING POISSON BRACKET $\dot{f} = \{f, \mathcal{H}\}$

LAGRANGIAN MECHANICS

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LAGRANGIAN MECHANICS

- „LIVES ON“ THE TANGENT BUNDLE TM
- LAGRANGIAN FUNCTION $\mathcal{L}: TM \rightarrow \mathbb{R}$ + VARIATIONAL CALCULUS GIVES EULER-LAGRANGE EQUATIONS AND DEFINITION OF MOMENTA

OR

- LAGRANGIAN FUNCTION + TULCZYJEW MAP $\alpha_M: TT^*M \rightarrow T^*TM$ GIVES PHASE EQUATIONS ON T^*M

QUESTION: WHAT GEOMETRIC STRUCTURE IS „RESPONSIBLE“ FOR LAGRANGIAN MECHANICS? \longrightarrow ALGEBROID

INTRODUCTION

HAMILTONIAN MECHANICS

LAGRANGIAN MECHANICS

- ASSOCIATED WITH A PHASE SPACE \mathcal{P} WHICH IS QUOTIENT $(\mathcal{P}_1 \times \mathcal{P}_2) / \mathbb{R}$

- „LIVES ON“ THE TANGENT BUNDLE TM

(\mathcal{P}_1, \dots)

WHY?

QUANTUM MECHANICS IS USUALLY FORMULATED USING A CONCEPT OF A HAMILTONIAN, BUT FIELD THEORIES - CLASSICAL, QUANTUM, \mathcal{T} -MODELS... ARE BASED ON LAGRANGIAN FUNCTIONS / FORMS

$TM \longrightarrow \mathbb{R}$

S EULER -
FUNCTION

• ω

VECT

FUNCT

SYMP

POISS

2YJEW MAP

IN DARBOUX COORDINATES

$$\dot{q}^i = \frac{\partial \mathcal{H}}{\partial p_i} \quad \dot{p}^j = - \frac{\partial \mathcal{H}}{\partial q^j}$$

GIVES PHASE EQUATIONS ON T^*M

USING POISSON BRACKET

$$\dot{f} = \{f, \mathcal{H}\}$$

QUESTION: WHAT GEOMETRIC STRUCTURE IS „RESPONSIBLE“ FOR LAGRANGIAN MECHANICS? \longrightarrow ALGEBROID



TULCZYJEW TRIPLE

M - CONFIGURATION MANIFOLD

TM - POSITIONS & VELOCITIES

T^*M - POSITIONS & MOMENTA - PHASE SPACE

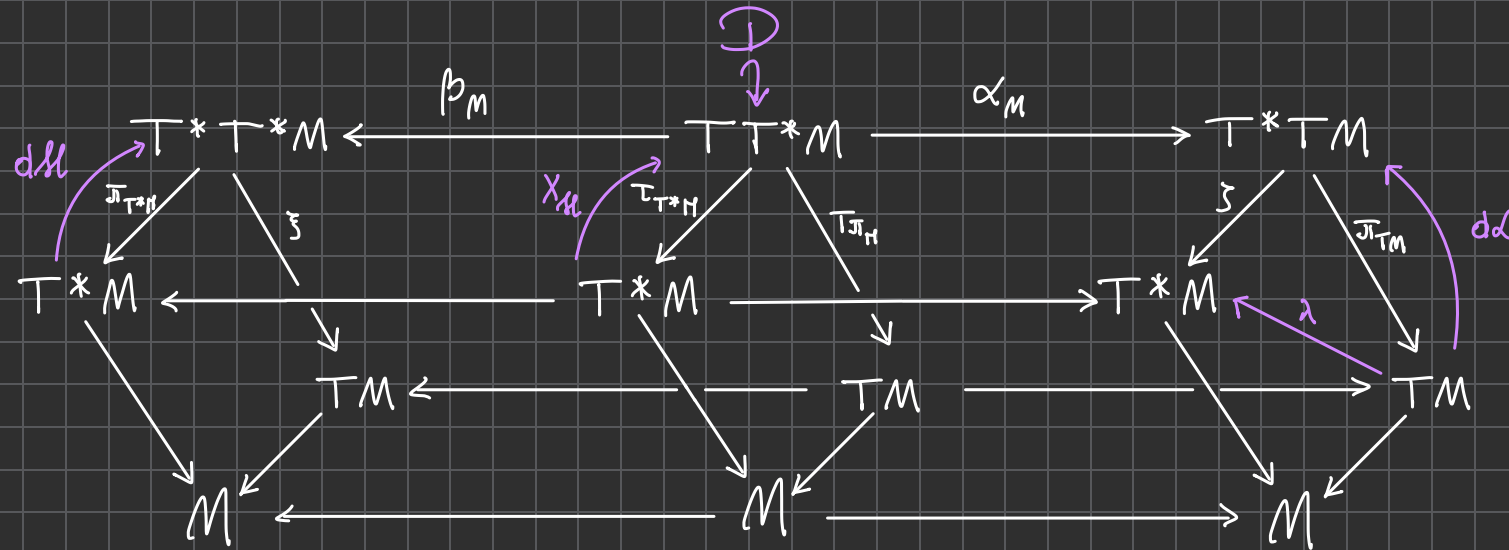


TULCZYJEW TRIPLE

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T^*M - POSITIONS & MOMENTA - PHASE SPACE



$$\mathbb{D} = X_M(T^*M) = \beta_M^{-1}(dL(T^*M))$$

$$\mathbb{D} = \alpha_M^{-1}(dL(TM))$$

BOTH CAN BE GENERALIZED

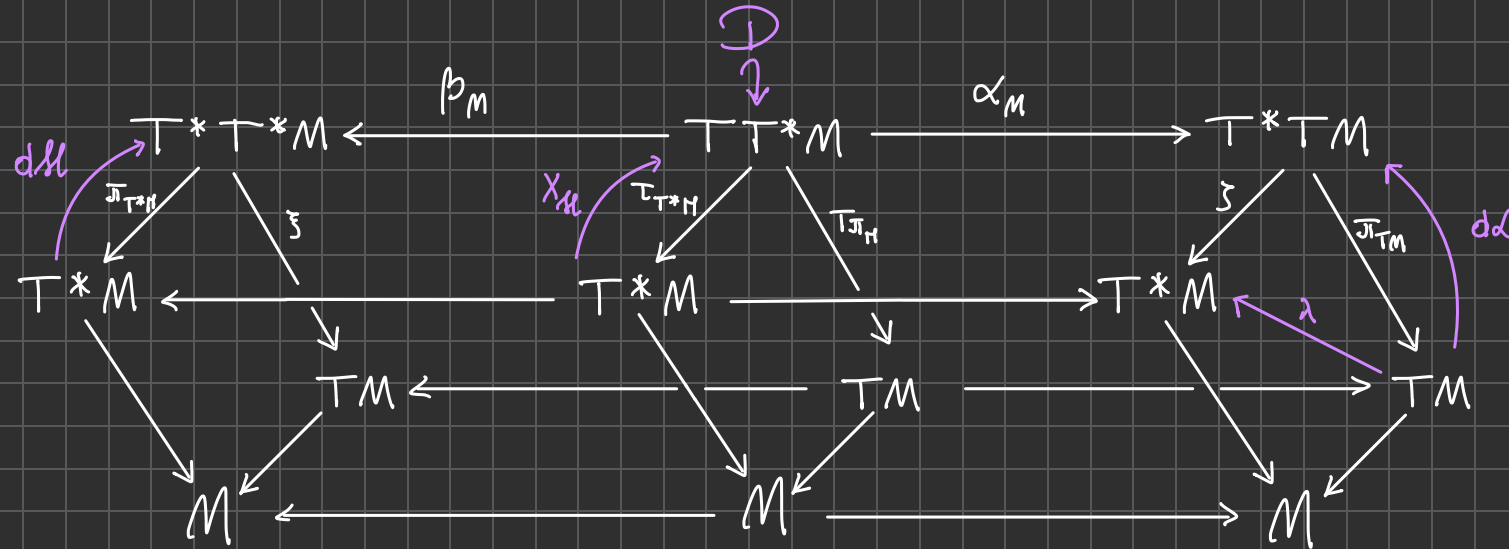


TULCZYJEW TRIPLE

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BOTH CAN BE GENERALIZED

$$\mathbb{D} = \beta_M^{-1}(S_{\mathcal{L}})$$

SOME LAGRANGIAN SUBMANIFOLD OF T^*T^*M

$$\mathbb{D} = \alpha_M^{-1}(S_{\mathcal{L}})$$

SOME LAGRANGIAN SUBMANIFOLD OF T^*TM

EXAMPLES - MASSIVE AND MASSLESS RELATIVISTIC PARTICLES

(+, -, -, -)
↓

IN THE FOLLOWING WE ASSUME THAT M IS AN AFFINE MINKOWSKI SPACE WITH CONSTANT METRIC η . ONE MAY WORK AS WELL WITH THE SPACE-TIME OF GENERAL RELATIVITY AS LONG AS WE TREAT PARTICLES AS TEST PARTICLES I.E. NOT BEING SOURCES OF THE GRAVITATIONAL FIELD.

$$M \ni q \quad TM = M \times V \quad T^*M = M \times V^* \quad \tilde{\eta}: V \rightarrow V^* \quad \tilde{\eta}(v) = \eta(v, \cdot) \quad \|v\| = \sqrt{\eta(v, v)}$$

$\eta(v, v) > 0$

$$T^*T^*M \xleftarrow{\beta_M} TT^*M \xrightarrow{\alpha_M} T^*TM$$

$$M \times V^* \times V^* \times V \quad M \times V^* \times V \times V^* \quad M \times V \times V^* \times V$$

$$(q, p, -\dot{p}, \dot{q}) \quad (q, p, \dot{q}, \dot{p}) \quad (q, \dot{q}, \dot{p}, p)$$

$$\mathcal{H}: M \times V^* \times \mathbb{R}_+ \rightarrow \mathbb{R}$$

$$\mathcal{H}(q, p, \tau) = \tau (\|p\|^2 - m^2)$$

$$\mathcal{D} \begin{matrix} \dot{q} \\ \dot{p} \end{matrix}$$

$$\left\{ (q, p, \tau \tilde{\eta}^{-1}(p), 0) : \langle p, \tilde{\eta}^{-1}(p) \rangle = m^2 \quad \tau > 0 \right\}$$

A FREE PARTICLE OF MASS m

$$L: M \times V_+ \rightarrow \mathbb{R}$$

$$L(q, \dot{q}) = m \|\dot{q}\|$$

$$S_L = \int dL(M \times V_+)$$

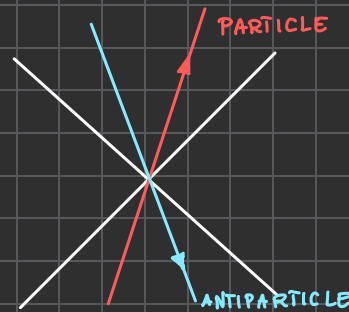
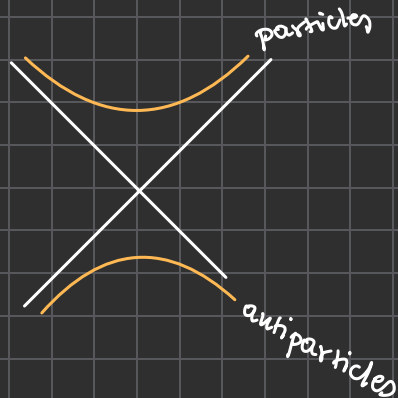
$$\left\{ (q, \dot{q}, 0, m \frac{\tilde{\eta}(\dot{q})}{\|\dot{q}\|}) \right\}$$

\dot{q} TIMELIKE

$$p = \lambda(\dot{q}) = m \frac{\tilde{\eta}(\dot{q})}{\|\dot{q}\|}$$

LEGENDRE MAP IS NOT INVERTIBLE

$$\frac{\partial \mathcal{H}}{\partial \dot{q}} = 0 \quad \parallel \dot{p}$$



GENERATING FAMILY:

$$F \xrightarrow{h} \mathbb{R}$$

$$\downarrow Q$$

$$(VF)^\circ \subset T^*F$$

$$\downarrow \mathcal{S}$$

$$T^*Q$$

$$\mathcal{D}_h = \mathcal{S} \left(dh \wedge (VF)^\circ \right)$$

+ CONDITIONS FOR h ENSURING
THAT \mathcal{D}_h IS ACTUALLY A SUBMANIFOLD

GENERATING FAMILY IN MECHANICS

$$\mathcal{H} : T^*M \times_n TM \longrightarrow \mathbb{R} \quad \mathcal{H}(q, \dot{q}, p) = \langle p, \dot{q} \rangle - L(q, \dot{q})$$

SOMETIMES GENERATING FAMILY CAN BE SIMPLIFIED

$$L(q, \dot{q}) = m \|\dot{q}\|$$

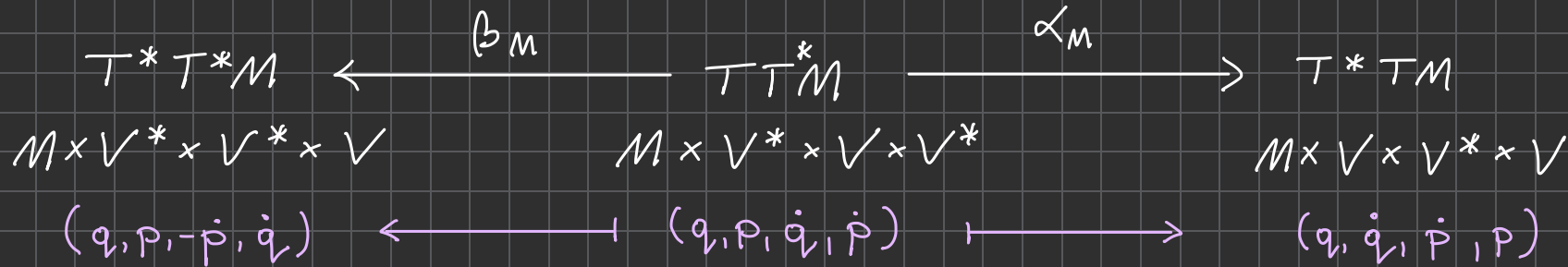
$$\mathcal{H}(q, p, \dot{q}) = \langle p, \dot{q} \rangle - L(q, \dot{q}) = \langle p, \dot{q} \rangle - m \|\dot{q}\|$$

$$\left\{ \begin{array}{l} \frac{\partial \mathcal{H}}{\partial \dot{q}} = 0 = p - \frac{m}{\|\dot{q}\|} \tilde{\eta}(\dot{q}) = p - m \tilde{\eta} \left(\frac{\dot{q}}{\|\dot{q}\|} \right) \Rightarrow \dot{q} = r \tilde{\eta}^{-1}(p) \end{array} \right.$$

VELOCITY IS PROPORTIONAL
TO $\tilde{\eta}^{-1}(p)$ WITH POSITIVE
COEFFICIENT!

$$\begin{aligned} \mathcal{H}(q, p, r) &= \langle p, r \tilde{\eta}^{-1}(p) \rangle - m \sqrt{r^2 \underbrace{\langle p, \tilde{\eta}^{-1}(p) \rangle}_{m^2}} = \\ &= r \|p\|^2 - r m^2 = \\ &= r (\|p\|^2 - m^2) \end{aligned}$$

MASSLESS PARTICLE



WORLD LINE OF A MASSLESS PARTICLE
LIES ON THE LIGHT CONE

$$\mathcal{D} = \{(q, p, \dot{q}, \dot{p}) : \|p\| = 0, \dot{p} = 0, \dot{q} = r \tilde{\eta}^{-1}(p)\}$$

$$H: M \times (V^*)^x \times \mathbb{R}_+ \longrightarrow \mathbb{R}$$

$$H(q, \dot{q}, y) = \frac{1}{2} y \gamma(p, p)$$

Non-zero VECTORS

$$L: M \times V^x \times \mathbb{R}_+ \longrightarrow \mathbb{R}$$

$$L(q, \dot{q}, y) = \frac{1}{2y} \gamma(\dot{q}, \dot{q})$$

THESE TWO GENERATING FAMILIES ARE
RELATED BY THE LEGENDRE TRANSFORMATION



W.M. Tulczyjew, P. Urbański

A Slow and Careful Legendre Transformation for Singular Lagrangians

vol. 30, p. 2909 (70 pages)

abstract • links/reference

Vol. 30 (1999), No. 10, pp. 2853 – 3027

ACTA PHYSICA
POLONICA **B**

LIE ALGEBROID

THERE EXISTS SEVERAL EQUIVALENT DEFINITIONS OF A LIE ALGEBROID ON A VECTOR BUNDLE $\tau: E \rightarrow M$

$$\textcircled{1} [\cdot, \cdot]: \text{Sec}(\tau) \times \text{Sec}(\tau) \longrightarrow \text{Sec}(\tau), \quad \begin{array}{ccc} E & \xrightarrow{\rho} & TM \\ \downarrow & & \downarrow \\ M & \xrightarrow{=} & M \end{array} \text{ VB MORPHISM}$$

ANTISYMMETRIC $[X, Y] = -[Y, X],$

JACOBI IDENTITY $[X, [Y, Z]] = [[X, Y], Z] + [Y, [X, Z]],$

$[X, fY] = f[X, Y] + \rho(X)(f)Y, \quad f \in C^\infty(M).$

EXAMPLES: $\left(\begin{array}{c} TM \\ \downarrow \\ M \end{array}, [\cdot, \cdot], \text{id}_{TM} \right)$

$\left(\begin{array}{c} \mathfrak{g} \\ \downarrow \\ \mathfrak{t} \end{array}, [\cdot, \cdot], 0 \right)$

FOR A PRINCIPAL BUNDLE $\begin{array}{c} P \ni G \\ \downarrow \\ M \end{array}$ WE DEFINE $AP = TP/G$

WITH SECTIONS BEING INVARIANT VECTOR FIELDS AND ρ COMING FROM $T\tau$

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$[X, fY] = f[X, Y] + \rho(X)(f)Y, \quad f \in C^\infty(M).$

$\textcircled{2}$ A LINEAR POISSON STRUCTURE Λ ON E^* ($\{\cdot, \cdot\}$)

COMPATIBLE WITH A VB STRUCTURE
HOMOGENEOUS OF DEGREE $-1: \mathcal{L}_{\nabla_{E^*}} \Lambda = -\Lambda$

CLOSED ON LINEAR FUNCTIONS

$\iota_x: E^* \rightarrow \mathbb{R}, \quad \iota_x(\varphi) = \langle \varphi, X \rangle$

$\{\iota_x, \iota_y\} = \iota_{[X, Y]}$

$\{\iota_x, \pi^* f\} = \pi^*(\rho(X)f)$

$\pi: E^* \rightarrow M$

① BRACKET OF $\text{Sec}(\tau)$ & ANCHOR

② LINEAR POISSON STRUCTURE ON E^*

③ A HOMOLOGICAL DERIVATION $d_E : \mathcal{A}(E^*) \longrightarrow \mathcal{A}(E^*)$ OF DEGREE 1
IN THE GRASSMANN ALGEBRA $\mathcal{A}(E^*)$ OF E^*

$$d_E : \mathcal{A}^i(E^*) \longrightarrow \mathcal{A}^{i+1}(E^*),$$

$$d_E^2 = 0,$$

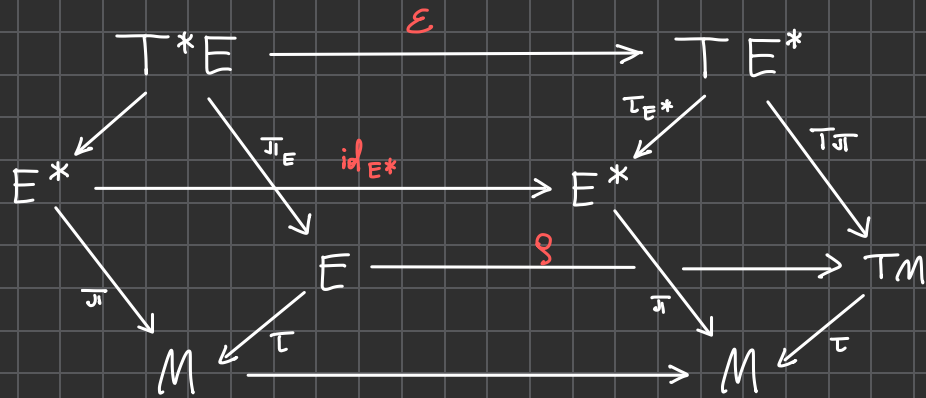
$$d_E(\alpha \wedge \beta) = (d_E \alpha) \wedge \beta + (-1)^a \alpha \wedge d_E(\beta), \quad \alpha \in \mathcal{A}^a(E^*).$$

RELATION WITH THE BRACKET

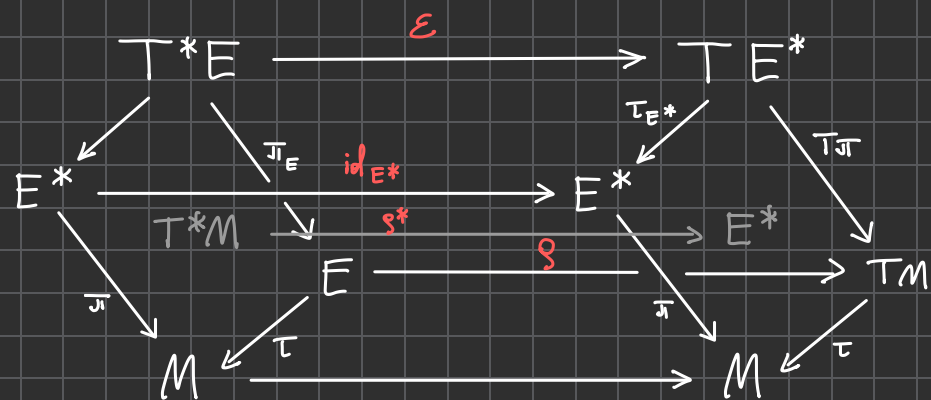
$$f \in C^\infty(M) = \mathcal{A}^0(E^*), \quad d_E f \in \mathcal{A}^1(E^*) \quad \langle d_E f, X \rangle = \rho(X)f,$$

$$\alpha \in \mathcal{A}^1(E^*) \quad d_E \alpha(X, Y) = \rho(X) \langle \alpha, Y \rangle - \rho(Y) \langle \alpha, X \rangle - \langle \alpha, [X, Y] \rangle.$$

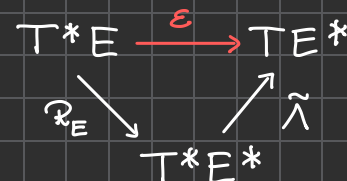
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- ② LINEAR POISSON STRUCTURE ON E^*
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- ④ A DVB MORPHISM $\varepsilon: T^*E \longrightarrow TE^*$ OVER THE IDENTITY ON E^*



- 1 BRACKET OF $\text{Sec}(\tau)$ & ANCHOR
- 2 LINEAR POISSON STRUCTURE ON E^*
- 3 A HOMOLOGICAL DERIVATION d_E OF DEGREE 1
- 4 A DVB MORPHISM $\varepsilon: T^*E \longrightarrow TE^*$ OVER THE IDENTITY ON E^*



RELATION WITH LINEAR POISSON STRUCTURE



COORDINATES

$$\begin{array}{ll}
 M : (q^i) & TE^* (q^i, \xi_a, \dot{q}^i, \dot{\xi}_b) \\
 E : (q^i, y^e) & T^*E (q^i, y^e, p_i, \eta_b) \\
 E^* : (q^i, \xi_a) &
 \end{array}$$

$$\begin{array}{l}
 q^i \circ \varepsilon = q^i \\
 \xi_a \circ \varepsilon = \eta_b \\
 \dot{q}^i \circ \varepsilon = g_a^i(q) y^e \\
 \dot{\xi}_b \circ \varepsilon = G_{db}^a(q) y^d \eta_c + g_c^i(q) p_i
 \end{array}$$

$$\varepsilon(q^i, y^e, p_i, \eta_b) = (q^i, \eta_b, g_a^i(q) y^e, G_{db}^a(q) y^d \eta_c + g_c^i(q) p_i)$$

- ① BRACKET OF $\text{Sec}(\tau)$ & ANCHOR
- ② LINEAR POISSON STRUCTURE ON E^*
- ③ A HOMOLOGICAL DERIVATION d_E OF DEGREE 1
- ④ A DVB MORPHISM $\varepsilon: T^*E \longrightarrow TE^*$
- ⑤ IN SUPERGEOMETRIC LANGUAGE :

A HOMOLOGICAL VECTOR FIELD ON THE SUPERMANIFOLD
 \overline{TE} ASSOCIATED TO THE V.B. $E \longrightarrow M$

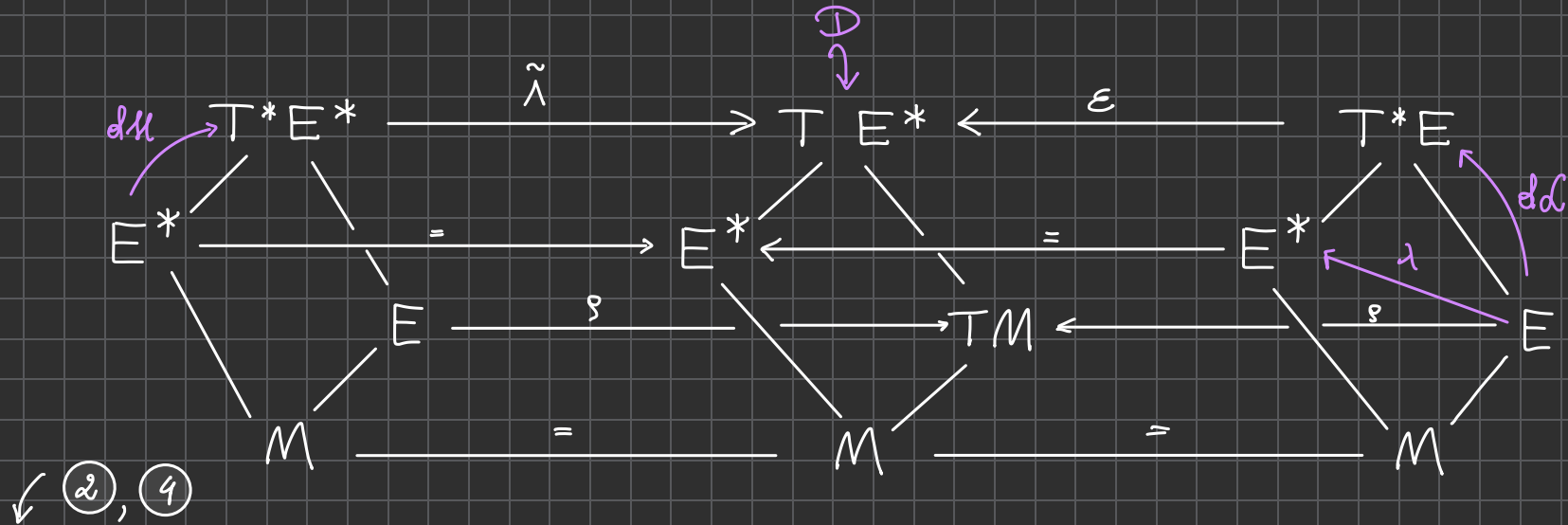
$$\mathcal{M} = \overline{TE}$$

$$C^\infty(\mathcal{M}) = \mathcal{A}(E^*)$$

VECTOR FIELD d_E



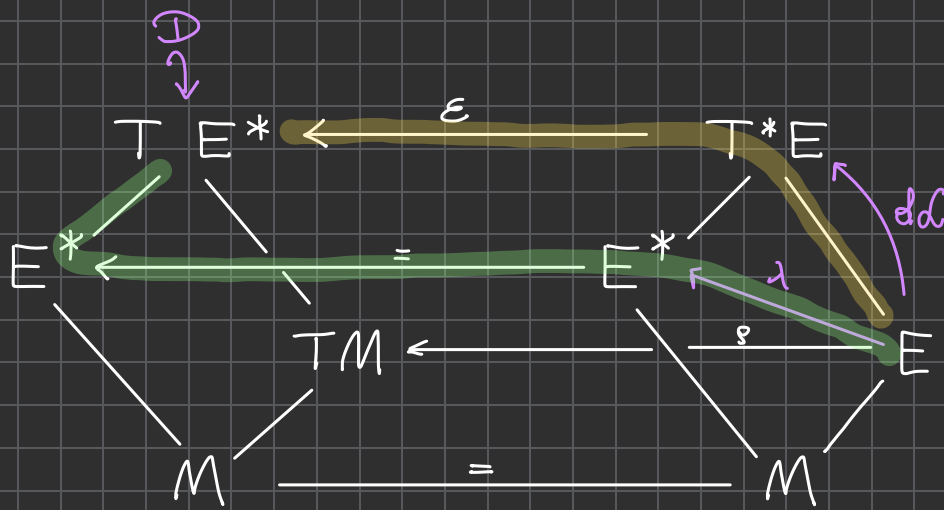
TULCZYJEW TRIPLE IN LIE ALGEBROID SETTING



THE TWO WAYS OF ENCODING A LIE ALGEBROID STRUCTURE PRODUCE THE TULCZYJEW TRIPLE. THE PHASE SPACE IS NOW E^* , WITH $\mathbb{D} \subset TE^*$.

\mathbb{D} CAN BE GENERATED BOTH FROM LAGRANGIAN AND HAMILTONIAN GENERATING OBJECT.

EULER-LAGRANGE EQUATIONS IN LIE ALGEBROID SETTING



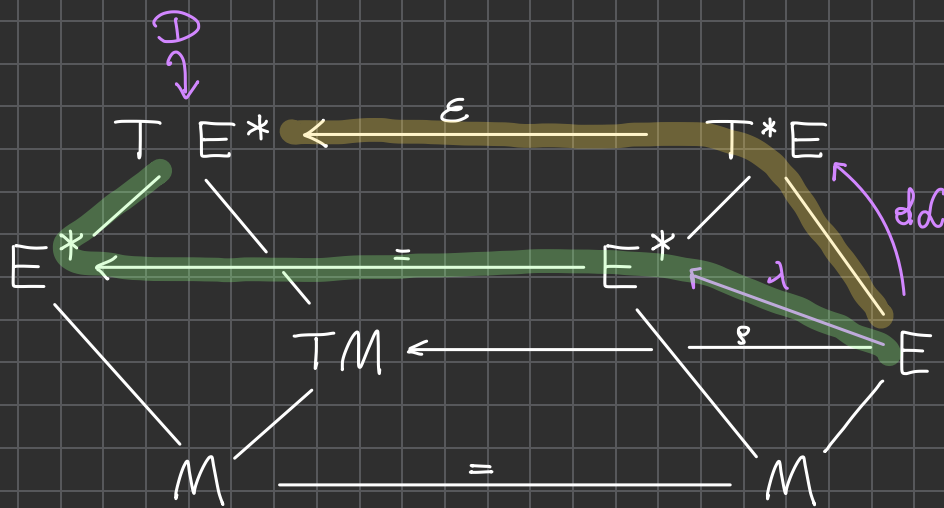
$$\mathcal{E}(q^i, y^a, p_i, \eta_b) = (q^i, \eta_b, g_a^i(q) y^a, C_{db}^a(q) y^d \eta_c + g_c^i(q) p_i)$$

$$\gamma: I \longrightarrow E$$

$$\mathcal{E}(\mathcal{D}\mathcal{L}(\gamma(t))) = \frac{d}{dt} \lambda(\gamma(t))$$

$$\mathcal{E} \circ \mathcal{D}\mathcal{L} = \Lambda_T$$

EULER-LAGRANGE EQUATIONS IN LIE ALGEBROID SETTING



$$\gamma: I \longrightarrow E$$

$$\varepsilon(d\lambda(\gamma(t))) = \frac{d}{dt} \lambda(\gamma(t))$$

$$\varepsilon(q^i, y^e, p_i, \eta_b) = (q^i, \eta_b, \rho_a^i(q) y^e, C_{db}^a(q) y^d \eta_c + \rho_c^i(q) p_i)$$

$$\frac{d}{dt} q^i(t) = \rho_a^i(q(t)) y^e(t) \quad \leftarrow \text{ADMISSIBILITY CONDITION}$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial y^e} (q(t), y(t)) \right) = C_{db}^a(q(t)) y^d(t) \frac{\partial \mathcal{L}}{\partial y^a} (q(t), y(t)) + \rho_b^i(q(t)) \frac{\partial \mathcal{L}}{\partial q^i} (q(t), y(t))$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial y^e} \right) - C_{db}^a y^d \frac{\partial \mathcal{L}}{\partial y^e} - \rho_b^i \frac{\partial \mathcal{L}}{\partial q^i} = 0$$



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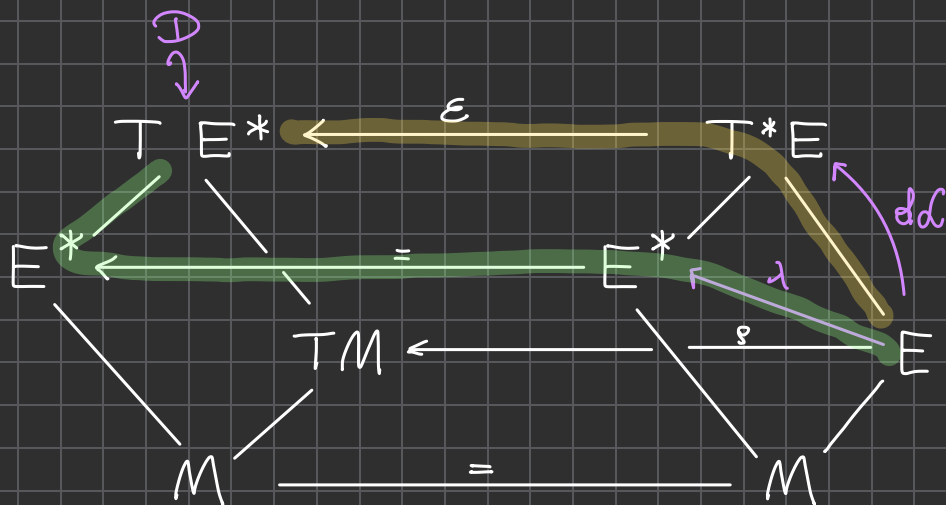
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GEOMETRICAL MECHANICS ON ALGEBROIDS

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EULER-LAGRANGE EQUATIONS IN LIE ALGEBROID SETTING



$$\gamma: I \longrightarrow E$$

$$\varepsilon(d\lambda(\gamma(t))) = \frac{d}{dt} \lambda(\gamma(t))$$

$$\varepsilon(q^i, y^e, p_i, \eta_b) = (q^i, \eta_b, \rho_a^i(q) y^e, G_{db}^a(q) y^d \eta_e + \rho_e^i(q) p_i)$$

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$$\frac{d}{dt} \left(\frac{\partial \lambda}{\partial y^e} (q(t), y(t)) \right) = G_{db}^a(q(t)) y^d(t) \frac{\partial \lambda}{\partial y^a} (q(t), y(t)) + \rho_b^i(q(t)) \frac{\partial \lambda}{\partial q^i} (q(t), y(t))$$

$$\frac{d}{dt} \left(\frac{\partial \lambda}{\partial y^e} \right) - G_{db}^a y^d \frac{\partial \lambda}{\partial y^e} - \rho_b^i \frac{\partial \lambda}{\partial q^i} = 0$$

EXERCISE FOR STUDENTS → WRITE EULER-LAGRANGE EQUATION $\frac{d}{dt} \left(\frac{\partial \lambda}{\partial \dot{q}^i} \right) - \frac{\partial \lambda}{\partial q^i} = 0$ IN COORDINATES (q^i, \dot{q}^i) ASSOCIATED

TO A BASIS OF SECTIONS OF $TM \rightarrow M$ DIFFERENT THAN $(\partial_{q^1}, \dots, \partial_{q^n})$. YOU WILL SEE SOMETHING LIKE THE E-L EQUATION ON LIE ALGEBROID WITH APPROPRIATE ρ AND G

DO WE SEE SITUATIONS LIKE THIS IN MECHANICS?

CHAIR OF MATHEMATICAL METHODS IN PHYSICS
KATEDRA METOD MATEMATYCZNYCH FIZYKI



MECHANICS ON LIE ALGEBROIDS - FIRST EXAMPLE

ALGEBROIDS IN MECHANICS APPEAR E.G. AS A RESULT OF REDUCTION WITH RESPECT TO SYMMETRIES. THE SIMPLEST EXAMPLE IS THE SYSTEM ON A LIE GROUP WITH INVARIANT LAGRANGIAN

G - A LIE GROUP

l_g - LEFT MULTIPLICATION
BY g

$$\mathcal{L}: TG \longrightarrow \mathbb{R} \quad \mathcal{L}(Tl_g(v)) = \mathcal{L}(v) \quad TG \cong G \times \mathfrak{g} \quad \mathcal{L}(g, x) = L(x)$$

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$$\mathcal{L}(g, x) = L(x)$$

FULL TULCZYJEW TRIPLE

$$T^*T^*G \xleftarrow{\beta_G} TT^*G \xrightarrow{\alpha_G} T^*TG$$

IN TRIVIALIZATION BY LEFT MULTIPLICATION

$$TG \cong G \times \mathfrak{g}, \quad T^*G \cong G \times \mathfrak{g}^*$$

$$G \times \mathfrak{g}^* \times \mathfrak{g}^* \times \mathfrak{g} \xleftarrow{\beta_G} G \times \mathfrak{g}^* \times \mathfrak{g} \times \mathfrak{g}^* \xrightarrow{\alpha_G} G \times \mathfrak{g} \times \mathfrak{g}^* \times \mathfrak{g}^*$$

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FULL TULCZYJEW TRIPLE

$$T^*T^*G \xleftarrow{\beta_G} TT^*G \xrightarrow{\alpha_G} T^*TG$$

IN TRIVIALIZATION BY LEFT MULTIPLICATION $TG \cong G \times \mathfrak{g}$, $T^*G \cong G \times \mathfrak{g}^*$

$$G \times \mathfrak{g}^* \times \mathfrak{g}^* \times \mathfrak{g} \xleftarrow{\beta_G} G \times \mathfrak{g}^* \times \mathfrak{g} \times \mathfrak{g}^* \xrightarrow{\alpha_G} G \times \mathfrak{g} \times \mathfrak{g}^* \times \mathfrak{g}^*$$

$$(g, \xi, \text{ad}_x^* \xi - \eta, X) \longleftarrow (g, \xi, X, \eta) \longrightarrow (g, X, \eta - \text{ad}_x^* \xi, \xi)$$



CALCULATED BY
MARCIN ZAJAC IN HIS
MASTER THESIS

IN TRIVIALIZATION $Tl_u(g, X) = (hg, X)$

INVARIANT LAGRANGIAN $\mathcal{L}(g, X) = \mathcal{L}(hg, X) = L(X)$

$$d\mathcal{L}(g, X) = (g, X, 0, dL(X))$$

$$G \times \mathfrak{g}^* \times \mathfrak{g}^* \times \mathfrak{g} \xleftarrow{\beta_G} \cancel{G} \times \mathfrak{g}^* \times \cancel{\mathfrak{g}} \times \mathfrak{g}^* \xrightarrow{\alpha_G} G \times \mathfrak{g} \times \mathfrak{g}^* \times \mathfrak{g}^*$$

$$K = G \times \mathfrak{g} \times \{0\} \times \mathfrak{g}^* \quad \text{COISOTROPIC SUBMANIFOLD}$$

TANGENT PROJECTION FROM GROUP ACTION

$$T\mathfrak{g}^*$$

$$\mathfrak{g}^* \times \mathfrak{g}^*$$

ε

$$\mathfrak{g} \times \mathfrak{g}^* = T^*\mathfrak{g}$$

SYMPLECTIC REDUCTION

AFTER REDUCTION I HAVE A MAP IN THE OTHER DIRECTION

$$G \times \mathfrak{g}^* \times \mathfrak{g}^* \times \mathfrak{g} \xleftarrow{\beta_G} \cancel{G} \times \mathfrak{g}^* \times \cancel{\mathfrak{g}} \times \mathfrak{g}^* \xrightarrow{\alpha_G} G \times \mathfrak{g} \times \mathfrak{g}^* \times \mathfrak{g}^*$$

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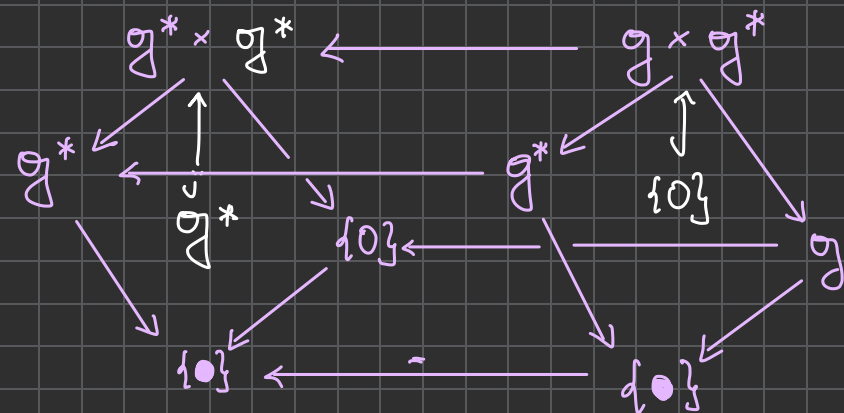
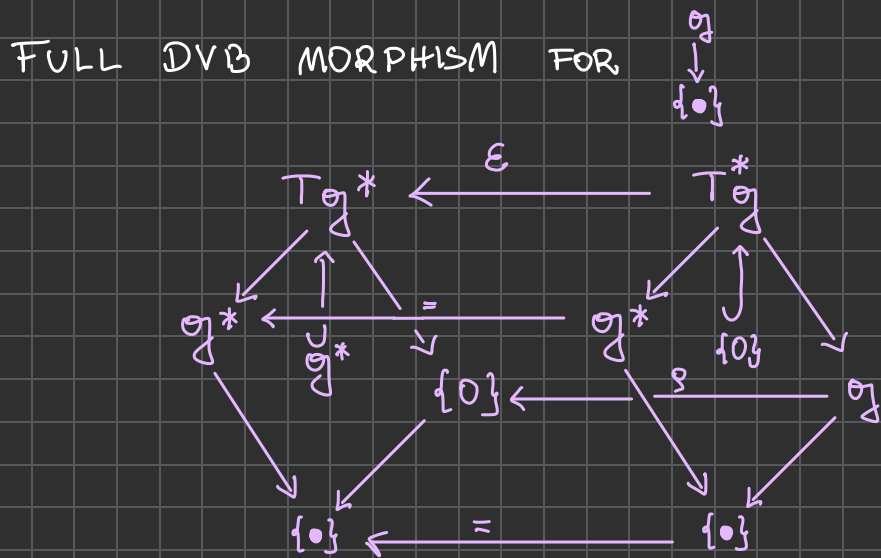
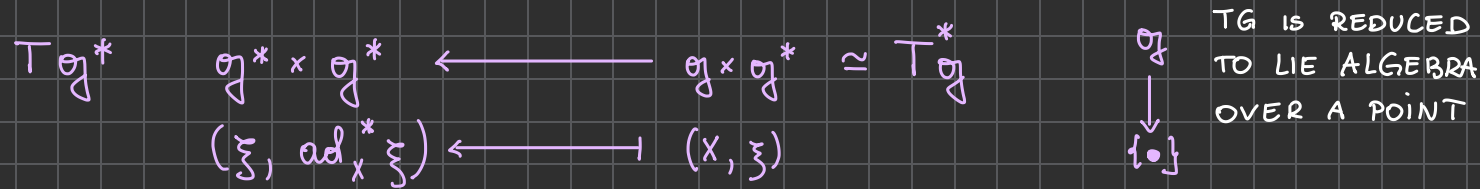
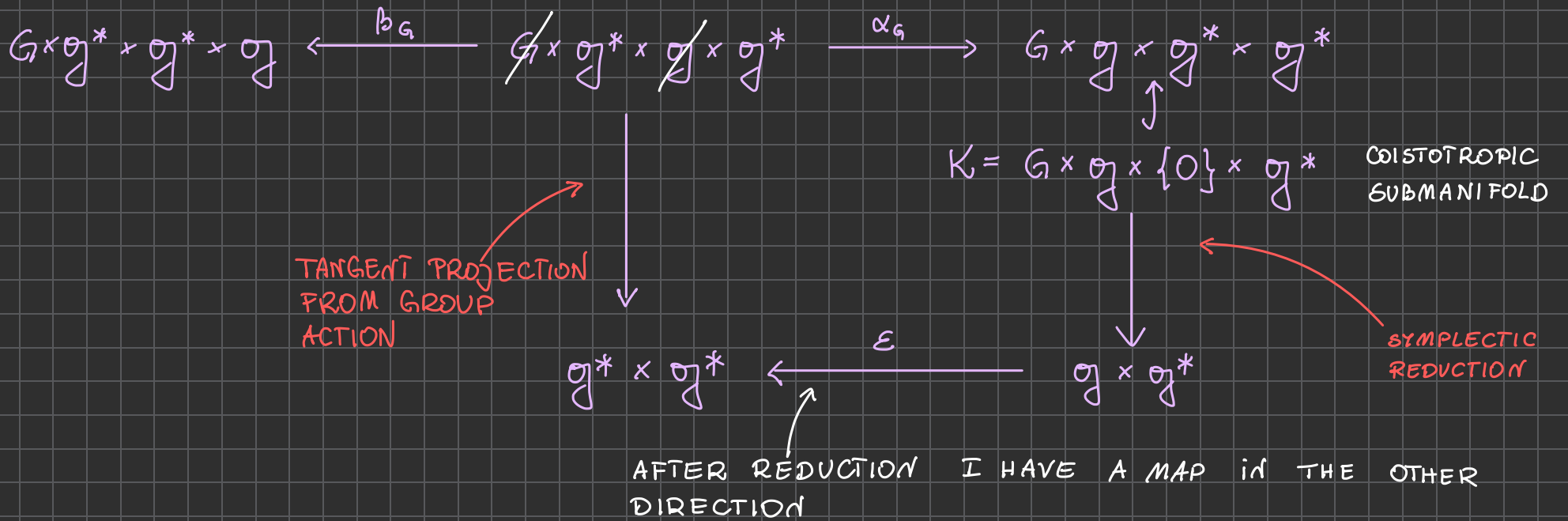
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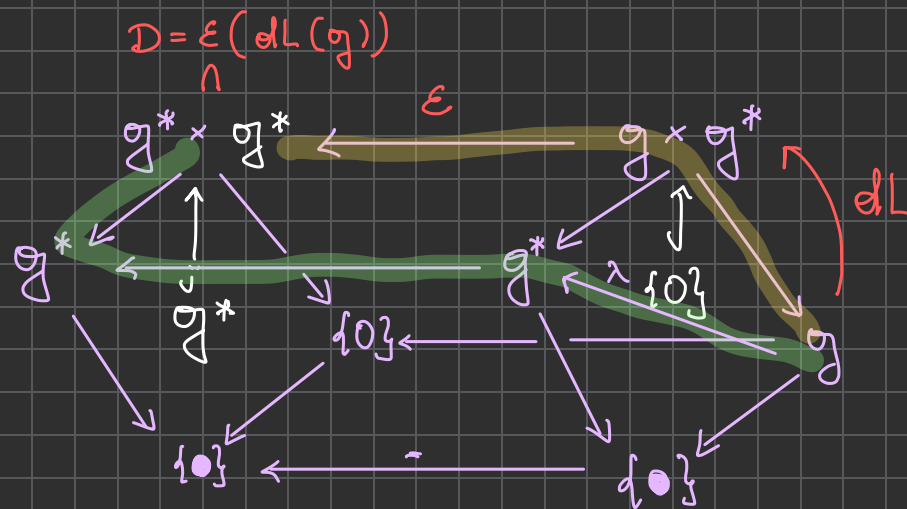
$$T\mathfrak{g}^* \quad \mathfrak{g}^* \times \mathfrak{g}^* \xleftarrow{\quad} \mathfrak{g} \times \mathfrak{g}^* \cong T\mathfrak{g}$$

$$(\xi, \text{ad}_x^* \xi) \xleftarrow{\quad} (x, \xi)$$

TG IS REDUCED TO LIE ALGEBRA OVER A POINT

$$\mathfrak{g} \downarrow \{ \bullet \}$$





LEGENDRE MAP: $\lambda: \mathfrak{g} \rightarrow \mathfrak{g}^* \quad x \mapsto dL(x)$

DYNAMICS $\mathcal{D} = \left\{ (\xi, \dot{\xi}) : \xi = dL(x) \quad \dot{\xi} = \text{ad}_x^*(dL(x)) \right\}$

EULER-Poincaré EQUATION ON \mathfrak{g} : $\frac{d}{dt} \left(\frac{\partial L}{\partial x} \right) = \text{ad}_x^* \frac{\partial L}{\partial x}$

EULER LAGRANGE ON \mathfrak{g}

$\mathfrak{g} \rightarrow \{\bullet\}$

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The Tulczyjew triple in mechanics on a Lie group

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FURTHER GENERALISATIONS : SKEW ALGEBROIDS ...

WE HAVE SEEN ON PREVIOUS SLIDES THAT WHAT WE REALLY USE IN MECHANICS IS A LIE ALGEBROID IN A FORM OF DOUBLE VECTOR BUNDLE MORPHISM. WE CAN THEN GENERALIZE IT DROPPING ASSUMPTIONS : JACOBI IDENTITY, ANTISYMMETRY ... UP TO JUST DVB MORPHISM $T^*E \longrightarrow TE^*$ OVER THE IDENTITY ON F^* **DO WE NEED IT?** SOME OF IT — MAYBE

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EXAMPLE : NONHOLONOMIC CONSTRAINTS + MECHANICAL LAGRANGIAN

WE START FROM LIE ALGEBROID (E, ε) AND SUBBUNDLE $K \subset E$ OVER M (FOR SIMPLICITY)

LAGRANGIAN IS ANY SMOOTH FUNCTION $\mathcal{L} : E \longrightarrow \mathbb{R}$

FROM VARIATIONAL APPROACH -
- d'ALEMBERT PRINCIPLE

(q^i, y^A, y^α) \swarrow CONSTRAINTS
 $K = \{y^\alpha = 0\}$

$$y^\alpha = 0$$

$$\dot{q}^i = \rho_A^i y^A$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial y^A} \right) - C_{DA}^B y^D \frac{\partial \mathcal{L}}{\partial y^B} - C_{DA}^B y^D \frac{\partial \mathcal{L}}{\partial y^B} - \rho_A^i \frac{\partial \mathcal{L}}{\partial q^i} = 0$$

THIS PART, PRESENT FOR ARBITRARY LAGRANGIAN, MEANS THAT WE CANNOT LOOK AT K ONLY

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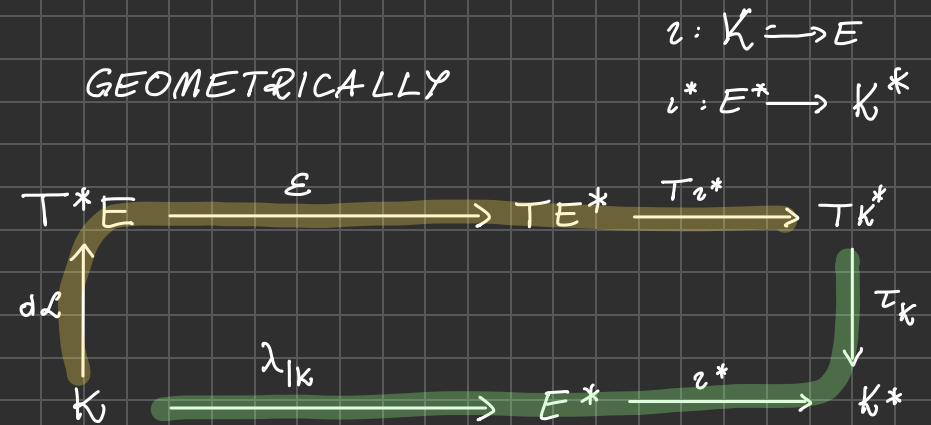
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FOR A MECHANICAL LAGRANGIAN:

g -BUNDLE METRIC ONE

$$g: E \times_m E \longrightarrow \mathbb{R}$$

BILINEAR
POSITIVE DEFINITE
SYMMETRIC

$$E = \mathcal{K} \oplus \mathcal{K}^\perp$$

$\uparrow \qquad \qquad \downarrow$
 $y^A \qquad \qquad y^\alpha$

$$g_{A\alpha} = 0$$

$$\mathcal{L}(q^i, y^A, y^\alpha) = \frac{1}{2} g_{AB} y^A y^B + \frac{1}{2} g_{\alpha\beta} y^\alpha y^\beta - V(q)$$

FOR A MECHANICAL LAGRANGIAN:

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 $g_{A\alpha} = 0$

$$\mathcal{L}(q^i, \dot{y}^A, \dot{y}^\alpha) = \frac{1}{2} g_{AB} \dot{y}^A \dot{y}^B + \frac{1}{2} g_{\alpha\beta} \dot{y}^\alpha \dot{y}^\beta - V(q)$$

THIS PART DOES NOT MATTER
I.E. WE CAN WORK WITH
LAGRANGIAN DEFINED
ON K ONLY

$$\left. \frac{\partial \mathcal{L}}{\partial y^\beta} = g_{\alpha\beta} \dot{y}^\alpha \right|_{y^\alpha=0} = 0$$

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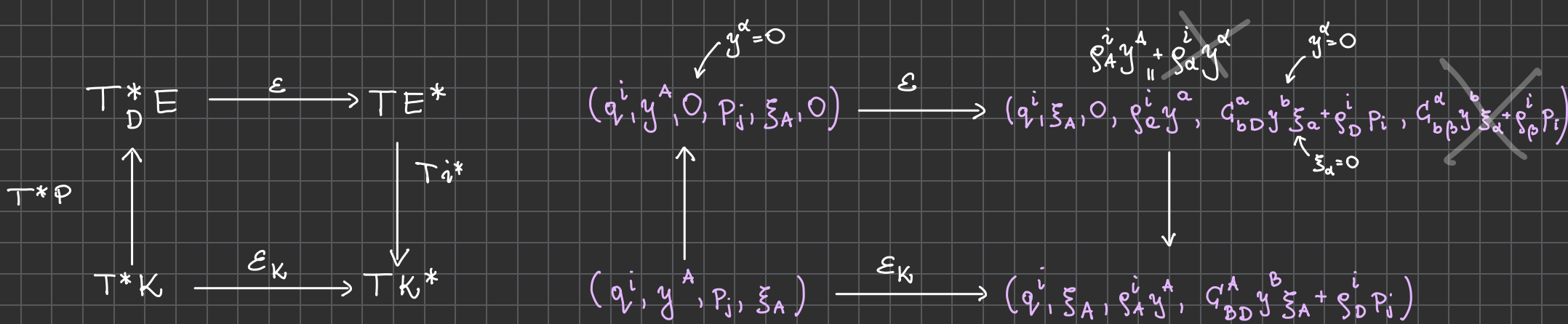
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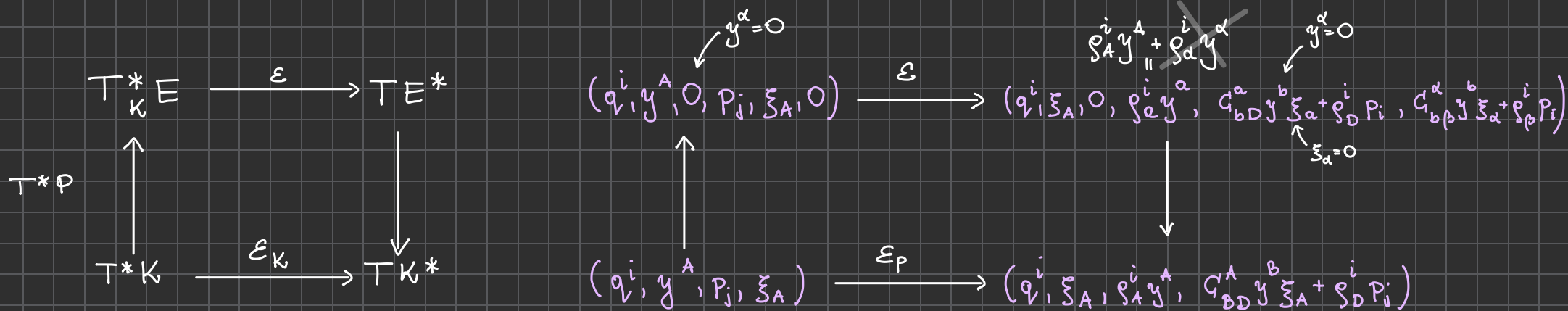
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$$E = K \oplus_M K^\perp \quad E^* = K^* \oplus_M K^0 \quad \iota: K \hookrightarrow E \quad \rho: E \rightarrow K \quad \alpha = (A, \alpha)$$



$$E = K \oplus_M K^{\perp} \quad E^* = K^* \oplus_M K^{\circ} \quad \iota: K \hookrightarrow E \quad \rho: E \longrightarrow K$$



WE HAVE A DOUBLE VECTOR BUNDLE MORPHISM $T^*K \longrightarrow TK^*$ OVER THE IDENTITY ON K^* CORRESPONDING TO A BIVECTOR ON K^* , BUT NO JACOBI IDENTITY

THE BRACKET ON K^* IS CALLED A **NONHOLONOMIC BRACKET**

THE THEORY WORKS FOR PARTICULAR LAGRANGIANS/HAMILTONIANS ONLY



Nonholonomic constraints: A new viewpoint

J. Grabowski; M. de León; J. C. Marrero; D. Martín de Diego

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WORKS FOR MECHANICAL LAGRANGIANS
WITH MAGNETIC FIELD-LIKE TERMS

DATA FROM THE LAGRANGIAN ENTER
THE STRUCTURE



Journal of Geometry and Physics
Volume 61, Issue 11, November 2011, Pages 2233-2253



Dirac algebroids in Lagrangian and Hamiltonian mechanics ☆

Katarzyna Grabowska^a ✉, Janusz Grabowski^b 👤 ✉

LINEAR CONSTRAINTS

NO RESTRICTIONS FOR LAGRANGIANS

REPLACE MAPS WITH RELATIONS

On Dirac Structures Admitting A Variational Approach

Oscar Cosserat, Alexei Kotov, Camille Laurent-Gengoux, Leonid Ryvkin,
Vladimir Salnikov

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CROCUSES IN TATRA MOUNTAINS PHOTO: PORTALTATRZANSKI.PL

THANK YOU FOR YOUR ATTENTION!